

AD/A-005 222

A FLUX CIRCUIT ANALYSIS FOR THE
MAGNETIC TRANSDUCER OF A FLUIDIC
REED GENERATOR

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Harry Diamond Laboratories

Prepared for:

Army Electronics Command

January 1975

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REPORT DOCUMENTATION PAGE		READ INSTRUCTIONS BEFORE COMPLETING FORM
1. REPORT NUMBER ECON-1,251	2. GOVT ACCESSION NO.	3. RECIPIENT'S CATALOG NUMBER ADJA 005 222
4. TITLE (and Subtitle) A FLUX CIRCUIT ANALYSIS FOR THE MAGNETIC TRANSDUCER OF A FLUIDIC REED GENERATOR		5. TYPE OF REPORT & PERIOD COVERED
7. AUTHOR(s) Herbert A. Leupold Frederick Rothwarf Doyle Edrington		6. PERFORMING ORG. REPORT NUMBER
8. CONTRACT OR GRANT NUMBER(s)		10. PROGRAM ELEMENT, PROJECT, TASK AREA & WORK UNIT NUMBERS Proj. No. 1W663613DE5506 Proj. No. A9-4-00068-01A9-CT
9. PERFORMING ORGANIZATION NAME AND ADDRESS Electronic Materials & Nuclear Hardening Rsch Area US Army Electronics Technology & Devices Lab (ECOM) Fort Monmouth, NJ 07703 AMSEL-TL-ES		12. REPORT DATE
11. CONTROLLING OFFICE NAME AND ADDRESS Harry Diamond Laboratories Connecticut Ave. at Van Ness NW Washington, DC 20438		13. NUMBER OF PAGES 55
14. MONITORING AGENCY NAME & ADDRESS (if different from Controlling Office)		15. SECURITY CLASS (of this report) Unclassified
16. DISTRIBUTION STATEMENT (of this Report) Approved for public release; distribution unlimited.		
17. DISTRIBUTION STATEMENT (of the abstract entered in Block 20, if different from Report) Approved for public release; distribution unlimited.		
18. SUPPLEMENTARY NOTES PRICES SUBJECT TO CHANGE		
19. KEY WORDS (Continue on reverse side if necessary and identify by block number) Artillery Fuzes Fluidic Generator Magnetic Permeances Eddy Current Losses		
20. ABSTRACT (Continue on reverse side if necessary and identify by block number) A mathematical analysis is made of the magnetic circuit of a fluidic generator to be used in various fuzing applications. A method for obtaining the approximate power output of a prototype generator as a function of the relevant geometric parameters is developed, and sample calculations are made for some typical values of these parameters. In the present design the calculations show that both the vibrating reed and part of the magnet keeper are magnetically saturated under projected operating conditions with the result that only about 35% of the potential power output is actually realized. (Cont'd on reverse side)		

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20. Abstract (cont'd)

Experimental data verify this conclusion, as agreement with the calculated value of power is better than 1%. Losses due to eddy currents and hysteresis are found to be much less serious amounting to no more than a few percent. Recommendations of possible remedies for the saturation problem are made. A computer study is made to determine the effect of variation of the relevant design parameters on the peak voltage output of the generator. The resulting design matrices provide a useful guide to design optimization as well as a clear delineation between the parameter combinations which result in saturation of the reed material and those which do not. Recommendations are made for design changes to improve the generator performance.

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A FLUX CIRCUIT ANALYSIS FOR THE MAGNETIC TRANSDUCER OF A FLUIDIC REED GENERATOR

INTRODUCTION

There is considerable interest in the utilization of fluidic generators in various fuzing applications.¹⁻³ These devices are mounted in the ogives of various projectiles and they essentially convert some of the energy of the relative motion of projectile and air first into acoustical energy and thence to electrical energy. A typical device is pictured schematically in Fig. 1.

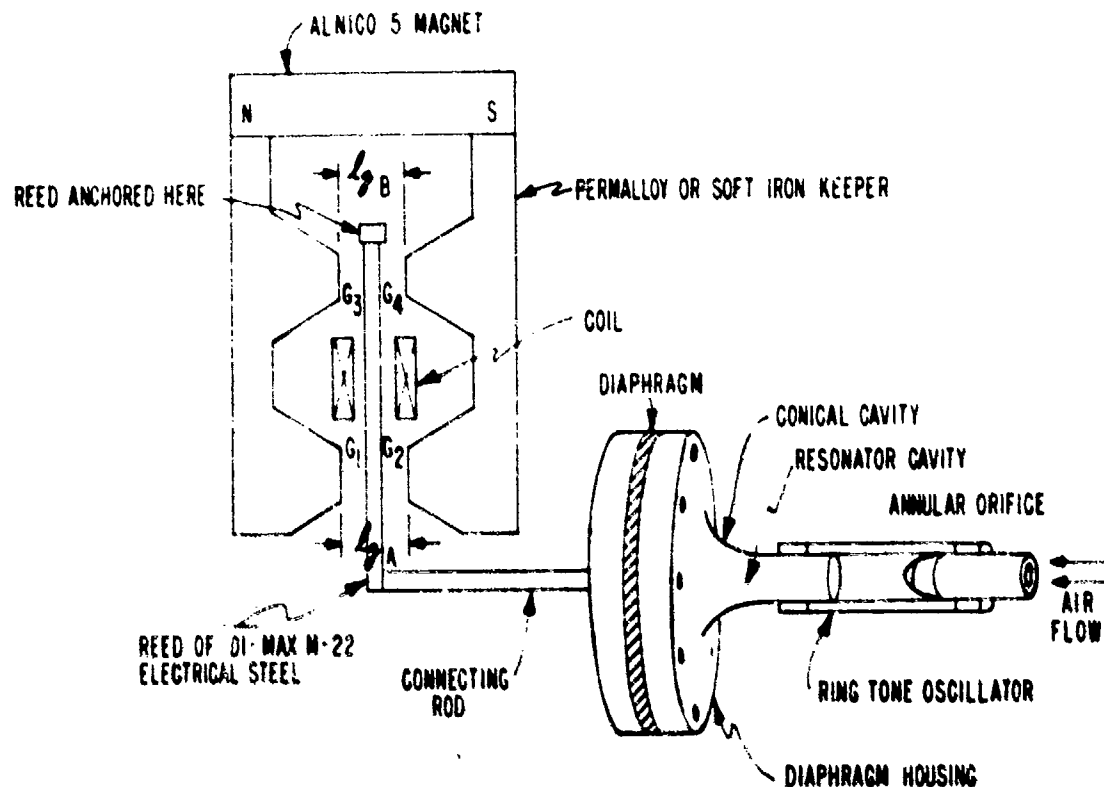


Fig. 1. Schematic View of Fluidic Generator with Reed Type Magnetic Transducer

1. Carl J. Campagnoulo, "The Fluidic Generator," HDL Technical Report-1328, September 1966.
2. Carl J. Campagnoulo, "Fluidic Power Generators for Ordnance Application," HDL Technical Report-1423, December 1968.
3. Carl J. Campagnoulo and Richard N. Gottron, "The Fluidic Generator: A New Electrical Power Source," 24th Annual Proceedings Power Sources Symposium, May 19-21, 1970.

As the projectile moves forward, air is forced through the annular orifice to excite the resonant cavity beyond. The oscillations of the cavity in turn cause the diaphragm to vibrate. This vibration is then transmitted to a permeable reed of Armco electrical steel. As can be seen from Fig. 1, when the reed is in its equilibrium position it is exactly midway between the pole pieces of both magnetic gaps and, therefore, carries no magnetic flux. The vibration of the reed, however, changes the relative lengths of gaps 1 and 2 and flux passes through the reed, alternating direction with the frequency of vibration. Gaps 3 and 4 are taken to remain constant and equal. The alternating flux through the reed induces an ac voltage in the coil around it and causes a current to flow in a circuit with a load resistance R_L . To optimize the design parameters we need an analysis of the magnetic circuit of the transducer. This will enable us to calculate the energy dissipated in R_L as a function of the gaps between pole pieces (l_g), the reed thickness (t), the reed displacement amplitude (a), and the frequency of vibration of the reed (f).

THE MAGNETIC CIRCUIT EQUATIONS

The magnetic circuit can be replaced by the analogous electric circuit as shown in Fig. 2, where the various electrical admittances have the values

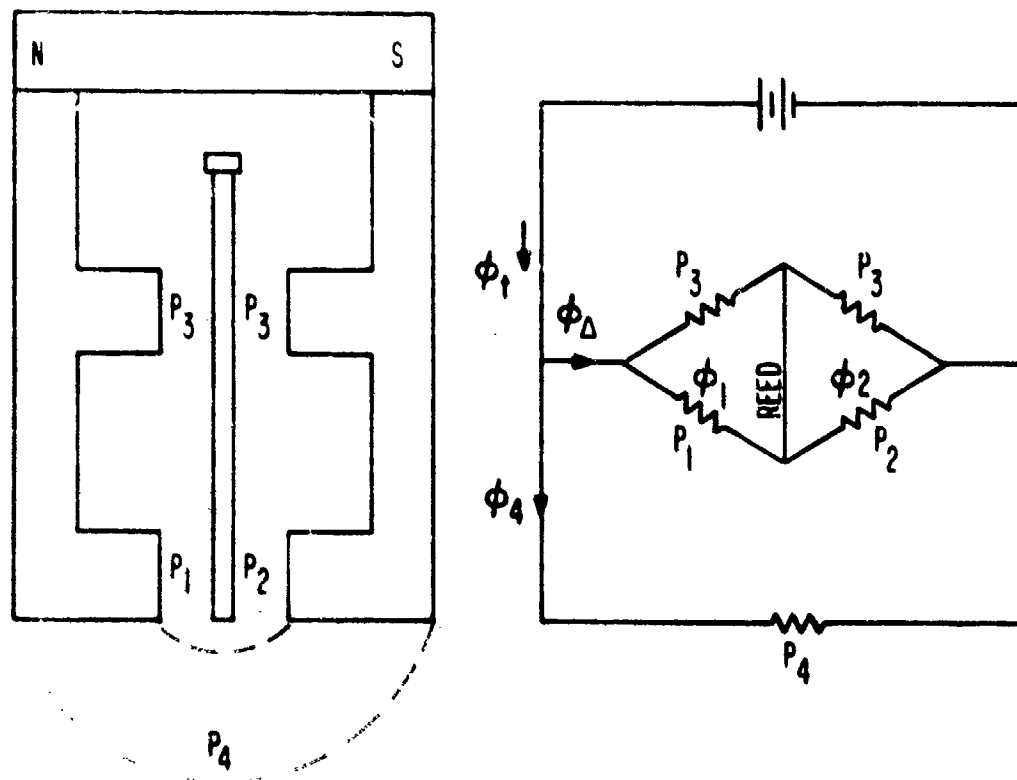


Fig. 2. Schematic Representation of the Magnetic Circuit and its Electrical Analogue

of the relevant magnetic permeances (P) of the reed generator, and the magnitudes of the currents deduced will then be the same as those of the fluxes (Φ) carried in the corresponding branches of the magnetic circuit. The circuit analysis is straightforward:

$$\Phi_{\Delta} = \frac{P_{\Delta}}{P_4 + P_{\Delta}} \Phi_t \quad (1)$$

$$\Phi_1 = \frac{P_1}{P_1 + P_3} \Phi_{\Delta} \quad (2)$$

$$\Phi_2 = \frac{P_2}{P_2 + P_3} \Phi_{\Delta} \quad (3)$$

$$\Phi_R = \Phi_1 - \Phi_2 \quad (4)$$

Substituting (1), (2), and (3) in (4) yields

$$\Phi_R = \left(\frac{P_1}{P_1 + P_3} - \frac{P_2}{P_2 + P_3} \right) \frac{\Phi_t}{1 + \frac{P_4}{P_{\Delta}}} \quad (5)$$

$$\frac{1}{P_{\Delta}} = \frac{1}{P_1 + P_3} + \frac{1}{P_2 + P_3} \quad (6)$$

$$P_{\Delta} = \frac{(P_1 + P_3)(P_2 + P_3)}{P_1 + P_2 + 2P_3} \quad (7)$$

Substitution of (7) in (5) finally gives us:

$$\Phi_R = \Phi_t \left(\frac{P_1}{P_1 + P_3} - \frac{P_2}{P_2 + P_3} \right) \frac{1}{1 + \frac{P_4(P_1 + P_2 + 2P_3)}{(P_1 + P_3)(P_2 + P_3)}} \quad (8)$$

$$\Phi_R = \Phi_t \left(\frac{P_1}{P_1 + P_3} - \frac{P_2}{P_2 + P_3} \right) \left(\frac{(P_1 + P_3)(P_2 + P_3)}{(P_1 + P_3)(P_2 + P_3) + P_4(P_1 + P_2 + 2P_3)} \right) \quad (9)$$

$$\Phi_R = \Phi_t \frac{P_1(P_2 + P_3) - P_2(P_1 + P_3)}{(P_1 + P_3)(P_2 + P_3) + P_4(P_1 + P_2 + 2P_3)} \quad (10)$$

$$\Phi_R = \Phi_t \frac{P_3(P_1 - P_2)}{(P_1 + P_3)(P_2 + P_3) + P_4(P_1 + 2P_3 + P_2)} \quad (11)$$

THE PERMEANCE CALCULATIONS

At this point it is necessary to obtain expressions for the various permeances P in terms of the reed displacement amplitude (a) and the configurational parameters which can be conveniently varied, namely the two gap lengths l_{gA} and l_{gB} , and the reed thickness t . Since l_{gA} is always kept equal to l_{gB} we denote both gap lengths by l_g .

These permeances will be calculated by essentially the same techniques used in Ref. 4. These consist principally of approximating the actual flux paths by circular arcs and then applying the permeance formulae for these paths as listed in the standard text by Roters.⁵ Suitable modifications of these general formulae must sometimes be made to deal with the peculiarities of a particular configuration. The actual arrangements for the prototype generator in both the isolated and in-shell conditions are shown in Figs. 3 and 4 respectively. The numerical values of the pertinent configurational dimensions are summarized in Table I.

The permeances associated with all of the flux paths which do not pass through the reed are collectively denoted by P_4 . There are five such paths which are schematically illustrated in Fig. 5, and are calculated immediately below along with the other permeances illustrated in Fig. 5. All calculated permeances are in inches.

(1) P_{p4} (Fig. 5-A)

The flux path here is in the form of a parallel plate magnetic capacitor and the permeance is simply given by

$$P_{p4} = \frac{\text{Plate Area}}{\text{Plate Separation}} = \frac{A_p}{S} = \frac{(.7)(.365)}{.395} = .647'' \quad (12)$$

P_{p4} will be slightly modified by the presence of the reed, but since the reed is very thin compared to the separation of the keeper plates this modification is very small and will be neglected.

(2) P_{y4} (Fig. 5-B)

This is a standard permeance path for which the formula is

$$.318 y \ln \left(1 + \frac{2w}{l} \right) \quad (13)$$

where y is the common perimeter length of the keeper edges linked by the flux path. In the fluidic generator $y = 2.2$ inches. w is the edge thickness .07" and l is the diameter of the smaller of the two semicircular boundaries of the flux path (.395"). Insertion of these values into Eq. (13) yields

4. H.A. Leupold, F. Rothwarf, C.G. Campagnoulo, J. Fine, and H. Lee, "Magnetic Circuit Design Studies for an Inductive Sensor," ECOM Technical Report-4158, October 1973.

5. H. Roters, Electromagnetic Devices, (John Wiley & Sons, Inc., 1941).

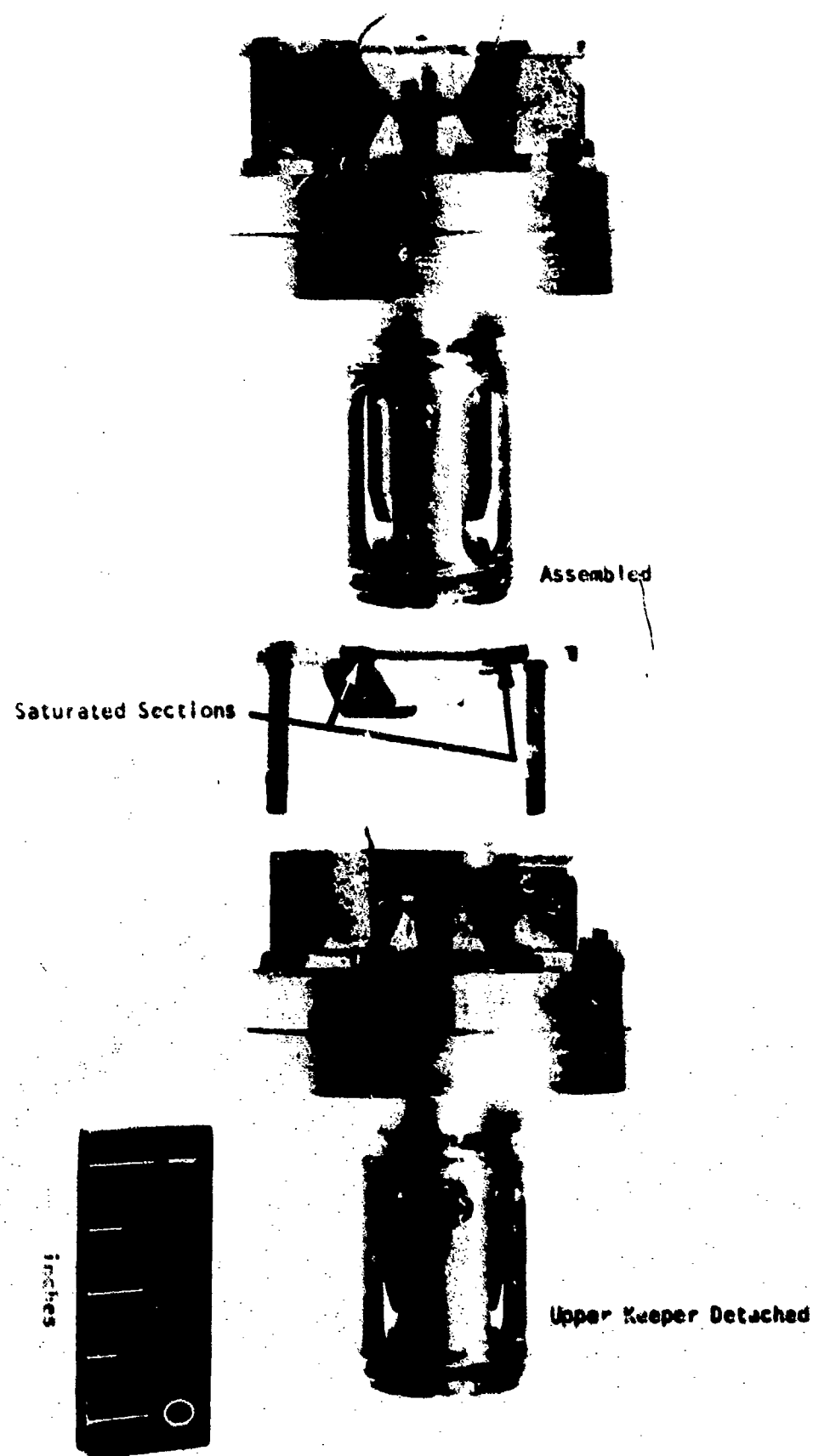
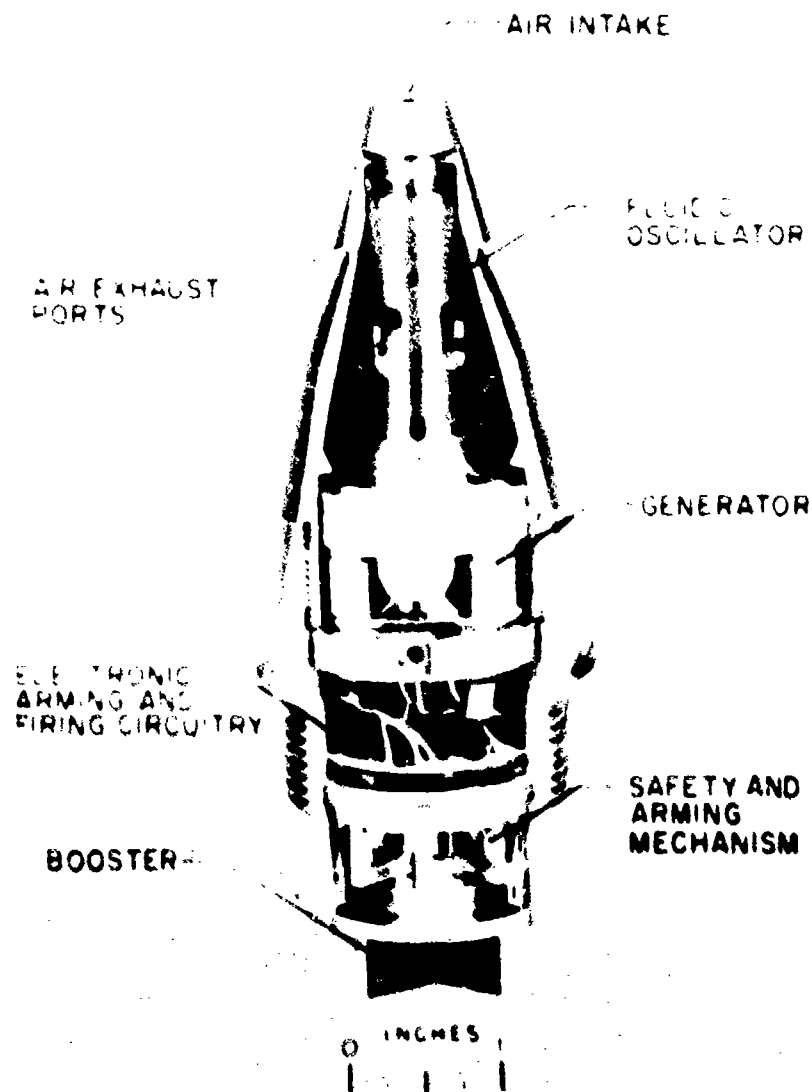


Fig. 3. Fluidic Generator.



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Fig. 4. Sectional View of Fluidic Generator
in Projectile Ogive

TABLE I

Fixed Dimensions of Generator Configuration and Values of Critical Constants

<u>Keeper</u>	<u>Inches</u>	<u>cm</u>	<u>Magnet</u>	<u>Inches</u>	<u>cm</u>
Plate area for P_{plate} calculation	0.256 in ²	1.65 cm ²	Length	0.395	1.00
Plate separation	0.395	1.00	Cross-section area	0.327 in ²	2.17 cm ²
Thickness	0.070	0.178	Effective outer leakage perimeter	2.3	5.8
Width (narrowest)	0.360	0.914			
Perimeter (effective)	2.2	5.6			
Saturation flux density		16,000 gauss			
Resistivity ρ		45 $\mu\Omega$ cm			
Hysteresis constant η		0.0001			
<u>Foles</u>	<u>Inches</u>	<u>cm</u>	<u>Reed</u>	<u>Inches</u>	<u>cm</u>
Face thickness	0.070	0.178	Width	0.270	0.686
Face width	0.313	0.795	Effective length	0.500	1.27
Face perimeter	0.766	1.95	Thickness *	0.0320	0.0812
Face area	0.0219 in ²	0.141 cm ²	Projection beyond gap	0.150	0.381
Average width	0.475	1.21	Vibrational amplitude *	0.017	0.043
Height (inner)	0.150	0.381	Volume *	0.00432 in ³	0.0706 cm ³
Height (outer)	0.235	0.597	Saturation flux density B_m	19,900 gauss	
Gap length *	0.070	0.178	Resistivity ρ	50 $\mu\Omega$ cm	
			Hysteresis constant η	.002	
			Projection beyond pole pieces	0.15	0.38

* Values are for sample calculations only. They are variable for computer calculated values in Table II (see Appendix II).

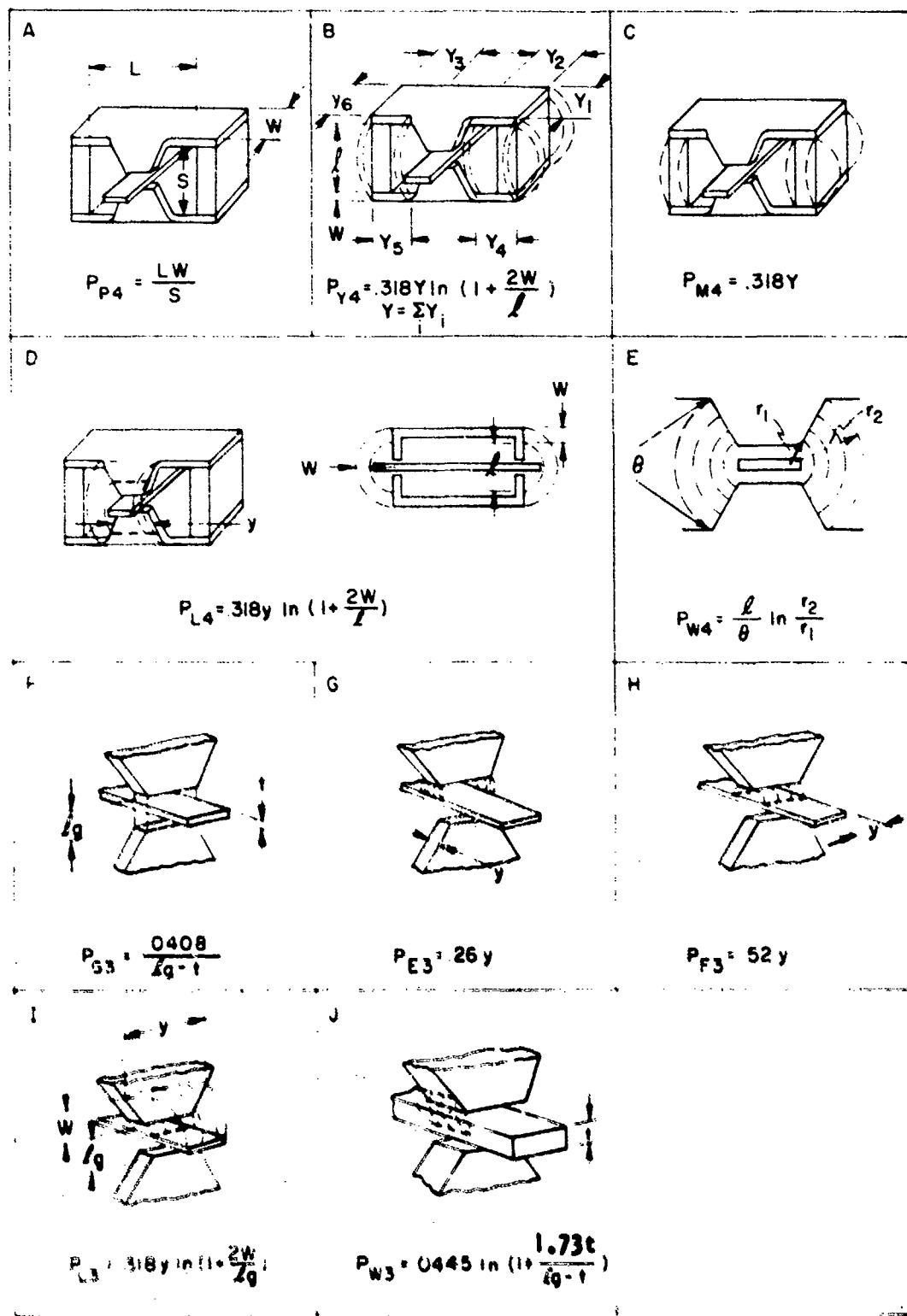


Fig. 5. Calculation of the Relevant Permeances

$$P_{y4} = .318 (2.2) \ln \left(1 + \frac{2(.07)}{.395} \right) = (.318)(2.2)(.303) = .212''$$

(3) P_{H4} (Fig. 5-C)

This is not a standard permeance but an expression for it can be derived as follows: Consider for the time being only that portion of the permeance path formed by the annular ring of width W as shown in Fig. 6. As we shall

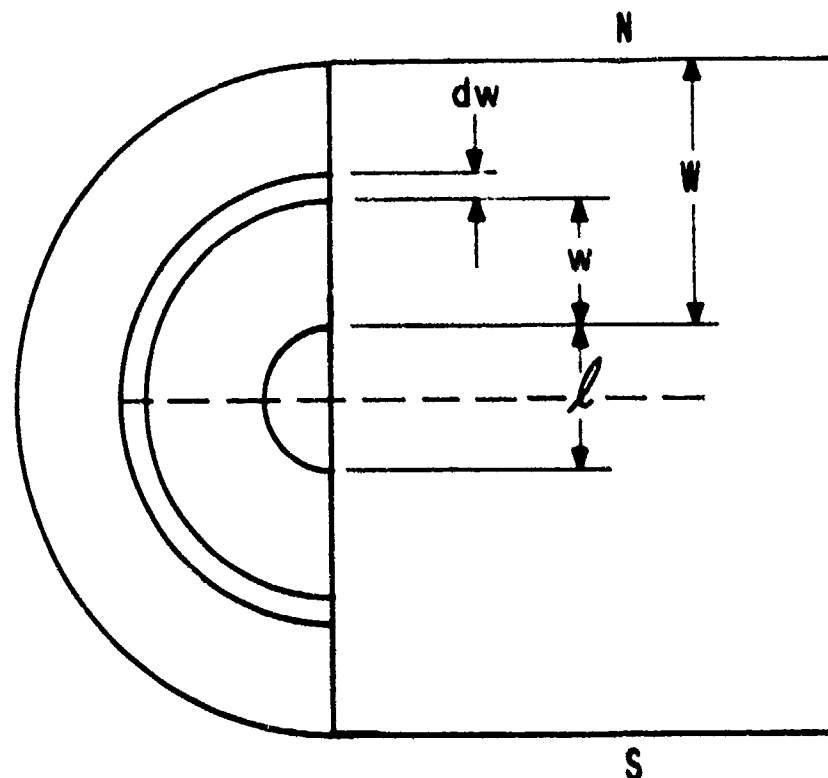


Fig. 6. Calculation of the Permeance P_{H4}

see a simple extension of the resulting formula can then be used to determine the remaining permeance consisting of the semicircular region of diameter d . The permeance of a path such as the annular ring is of the same standard form as Eq. (13), i.e.

$$P = .318 y \ln \left(1 + \frac{2W}{d} \right) \quad (14)$$

Similarly, the portion of the permeance within the annular ring of width w is

$$P = .318 y \ln \left(1 + \frac{2w}{d} \right) \quad (15)$$

and the permeance of element dw is

$$dP = \frac{.318y \frac{2}{l} dw}{1 + \frac{2w}{l}} \quad (16)$$

However, since the flux lines are emanating from the sides of a magnet, the magnetomotive force driving an element such as dw is only a fraction of the full MMF of the magnet. Since this MMF is in direct proportion to the distance of the path terminals from the magnet center, this fraction is equal to

$$f = \frac{\frac{l}{2} + w}{\frac{l}{2} + W} \quad (17)$$

and the effective permeance of dP is reduced by the factor f , viz.,

$$dP_{\text{eff}} = fdP = .318y \frac{\left(\frac{l}{2} + w\right) \frac{2}{l}}{\left(\frac{l}{2} + W\right) \left(1 + \frac{2w}{l}\right)} dw$$

$$dP_{\text{eff}} = \frac{.318y}{\left(\frac{l}{2} + W\right)} \quad (18)$$

and the total permeance of the annular ring W is

$$P_{\text{eff}} = \frac{.318y}{W + \frac{l}{2}} \int_0^W dw = \frac{.318y W}{W + \frac{l}{2}} \quad (19)$$

To obtain the total effective permeance P_{M4} , one needs only to set $l = 0$ in the expression for P_{eff} and we are left with

$$P_{M4} = .318y \quad (20)$$

y is just the perimeter of the magnet minus the length of the inner wall. The latter length is subtracted because it is rendered largely ineffective by competition from the flux paths of permeance P_{p4} . The effective y is somewhat awkward to estimate because of irregularities such as the screw slots cut into the magnet walls, but 2.3" seems a reasonable value for the two magnets together. Accordingly, our value for P_{M4} is

$$P_{M4} = (.318)(2.3) = .731"$$

(4) P_{L4} (Fig. 5-D)

This again is a standard permeance of the form used in P_{y4} . There is one such path for the outside surface of each pole piece so we must multiply Eq. (13) by two. The flux lines emanating from the inner surfaces

go to the reed and hence contribute to P_1 , P_2 , and P_3 rather than to P_{L4} . For y we take the average width, (.475") of the effective portion of the trapezoidal outer surface of each pole piece. L , (.30") is twice the length of the reed projection beyond the gap as is clear from Fig. 5-D. As .235" is the height of the trapezoidal pole piece plus the gap from the pole face to reed, it also follows from Fig. 5-D that $w = (.235 - .150) = .085$ ". Inserting these values in Eq. (13) yields

$$P_{L4} = 2(.318)(.475) \ln \left(1 + \frac{2(.085)}{.300} \right)$$

$$P_{L4} = .302 \ln 1.57 = .136"$$

(5) P_{w4} (Fig. 5-E)

The formula for a path of this kind is

$$P = \frac{L}{\theta} \ln \frac{r_2}{r_1} \quad (21)$$

and since we have four such paths (2 for each pair of pole pieces)

$$P_{w4} = 4 \frac{L}{\theta} \ln \frac{r_2}{r_1} \quad (22)$$

L is just the thickness of our pole pieces which is .070" and θ is $\frac{2}{3} \pi$ radians. r_2 is the distance from the vertex of angle θ to the larger base of the trapezoidal pole piece. Thus, r_2 is equal to the length of the pole piece edge plus one-half the gap width divided by the sine of $\frac{\pi}{3}$ radians. So since the pole piece edge is .2" long

$$r_2 = .2 + \frac{L_g}{2} \csc \frac{\pi}{3} = .2 + .577 L_g \quad (23)$$

r_1 is equal to $\frac{L_g}{2} \csc \frac{\pi}{3}$ plus the distance along the pole piece edge from which the flux lines go to the edge of the reed. The latter distance is just one-half the reed thickness (t). So finally we have

$$r_1 = \frac{t}{2} + .577 L_g \quad (24)$$

and

$$P_{w4} = \frac{(4)(.070)(3)}{2\pi} \ln \frac{.2 + .577 L_g}{\frac{t}{2} + .577 L_g} \quad (25)$$

$$P_{w4} = .134 \ln \frac{.4 + 1.16 L_g}{t + 1.16 L_g} \quad (26)$$

(6) P_4

We now have expressions for all of the components of the "leakage permeance" P_4 , so P_4 is given by:

$$P_4 = P_{p4} + P_{y4} + P_{M4} + P_{L4} + P_{w4}$$

$$P_4 = .647 + .212 + .731 + .136 + .134 \ln \frac{.4 + 1.16 \ell_g}{t + 1.16 \ell_g}$$

$$P_4 = 1.72 + .134 \ln \frac{.4 + 1.16 \ell_g}{t + 1.16 \ell_g} \quad (27)$$

If the gap is varied in the usual manner, i.e. by placing permalloy shims between the magnets and the permalloy keeper, P_{w4} , P_{p4} , and P_{y4} will be slightly affected, but since the shim thicknesses are only of the order of 10 or 20 thousandths of an inch the change in these quantities will be negligible.

(7) P_{G3} (Fig. 5-f)

This is the permeance of the gap between the reed and either pole piece at the pinned end of the reed. The length of this gap L is given by

$$L = \frac{\ell_g - t}{2} \quad (28)$$

Since the widths of the pole piece (.313") and the reed .270" are not the same, we take an average (.291") in our determination of the gap's cross-sectional area. The pole face thickness is .070" so

$$P_{G3} = \frac{A_{rcs}}{L} = \frac{2(.070)(.291)}{\ell_g - t} \quad (29)$$

$$P_{G3} = \frac{.0408}{\ell_g - t} \quad (30)$$

If the gap sizes are varied by machining the pole faces the constant in this formula will change due to the alteration in cross-sectional area of the pole faces.

(8) P_{E3} (Fig. 5-G)

P_{E3} is the permeance of the path extending from the edges of the reed to the pole face edges parallel to the former. The standard expression for this type of permeance path is

$$P = .26y \quad (31)$$

where y is the edge length which here is the thickness of the pole piece (.070") so

$$P = (.26)(.070) = .0182" \quad (32)$$

Since we have two such edges the total permeance is

$$P_{E3} = 2P = (2)(.0182) = .0364" \quad (33)$$

(9) P_{F3} (Fig. 5-H)

These are the paths from the edges of the pole pieces to the flat portion of the reed. These are standard forms whose permeances are given by

$$P = .52y = .52(.291) \quad (34)$$

Since there are two paths we have

$$P_{F3} = 2P = 2(.52)(.291) = .303" \quad (35)$$

(10) P_{L3} (Fig. 5-I)

These permeances involve the paths between the walls of the pole pieces and the wide faces of the reed. The general expressions are the same as those for P_{L4} except for a factor of 2 to take into account that the flux lines here are quarter circles rather than semicircles as in P_{L4} . The gap length l used in the formula must be taken as twice the separation between the reed and the pole piece, i.e.,

$$l = 2 \left(\frac{l_g - t}{2} \right) = l_g - t.$$

The quantity w_i for the inner flux path is just the altitude of the trapezoidal pole piece which is .15". For the effective y one chooses an average of the reed width (.270") and the width of the trapezoid midway between its bases (.406"). This results in a value $y = .338"$.

The outer path must be treated slightly differently because the reed projects only .15" beyond the gap edge. Accordingly, w_o is just taken as

$$w_o = .15 - \frac{(l_g - t)}{2} \quad (36)$$

and the average value y is also slightly altered.

As is seen in Fig. 7 the flux lines from the reed extend only up to line AB which is .15" above the top surface of the reed. The average width of the flux-emanating region of the pole face is then

$$y_F = \frac{AB + .313}{2} \quad (37)$$

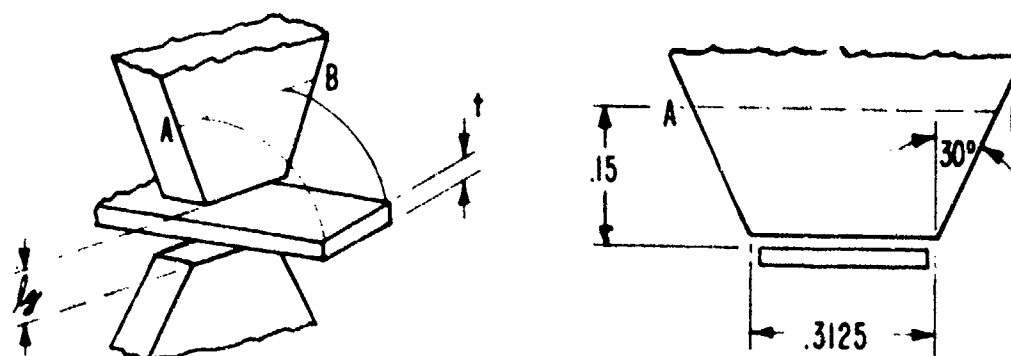


Fig. 7. Determination of the Width y_F in the Formula for Permeance P_{L3}

and

$$y_0 = \frac{y_F + y_R}{2} \quad (38)$$

where $y_R = .270$ is the reed width.

$$AB = .313 + 2 \left(.15 - \left(\frac{l_g - t}{2} \right) \right) \tan 30^\circ \quad (39)$$

$$AB = .313 + (.30 - (l_g - t)) .577$$

$$AB = .486 - .577 (l_g - t) \quad (40)$$

$$y_F = \frac{.486 - .577 (l_g - t) + .313}{2} = .399 - .289 (l_g - t) \quad (41)$$

$$y_0 = \frac{.399 - .289 (l_g - t) + .270}{2} \quad (42)$$

$$y_0 = .335 - .144 (l_g - t) \quad (43)$$

Substituting these values in the general expression yields

$$P_{L3} = 2(.318)(.338) \ln \left(1 + \frac{2(.15)}{l_g - t} \right) + 2(.318)[.335 - .144 (l_g - t)] \ln \left(1 + \frac{2 \left(.15 - \frac{l_g - t}{2} \right)}{l_g - t} \right) \quad (44)$$

$$P_{L3} = .215 \ln \left(1 + \frac{.30}{l_g - t} \right) + (.213 - .0916 (l_g - t)) \ln \left(1 + \frac{.30 - (l_g - t)}{l_g - t} \right) \quad (45)$$

$$P_{L3} = .215 \ln \left(1 + \frac{.30}{l_g - t} \right) + (.213 - .0916 (l_g - t)) \ln \left(\frac{.30}{l_g - t} \right) \quad (46)$$

(11) P_{w3} (Fig. 5-J)

If the narrow side face of the reed were coplanar with the narrow face of the pie piece as indicated by the dashed line in Fig. 5-J, the expression for permeance P_{w3} would be the often used standard form

$$P_{w3} = .318y \ln \left(1 + \frac{2w}{l} \right) \quad (47)$$

Although it is possible to correct for the fact that the actual flux lines are shorter than in this ideal case, the correction would be very small, and since the total permeance P_{w3} is itself so small as to barely affect the calculations with one significant figure, this correction is completely

negligible and will not be made here. So w is taken as $\frac{t}{2}$, $l = \frac{l_g - t}{2 \cos 30^\circ}$.

$y = .070$, and we precede the expression with the factor 2 because there are two paths of this type

$$P_{w3} = 2(.318)(.070) \ln \left(1 + \frac{2t \cos 30^\circ}{l_g - t} \right) \quad (48)$$

$$P_{w3} = .0445 \ln \left(1 + \frac{1.73t}{l_g - t} \right) \quad (49)$$

(12) P_{G1} and P_{G2}

These permeances are of the same form as P_{G3} with allowance made for the spacing changes due to the displacement of the reed. When the reed is at equilibrium, that is, at zero displacement, $P_{G1} = P_{G2} = P_{G3}$. Since we are interested in the maximum and minimum values of P_{G1} and P_{G2} , we write them in terms of the displacement amplitude (a) of the reed.

$$P_{G1}^{max} = P_{G2}^{max} = \frac{.0408}{l_g - t - 2a} \quad (50)$$

$$P_{G1}^{min} = P_{G2}^{min} = \frac{.0408}{l_g - t + 2a} \quad (51)$$

where P_{G1} is a maximum when P_{G2} is a minimum and vice versa.

(13) $\underline{P_{E1}}$

$$P_{E1} = P_{E3} = .0364'' \quad (52)$$

(14) $\underline{P_{F1}}$

$$P_{F1} = P_{F3} = .303'' \quad (53)$$

(15) $\underline{P_{W1} \text{ and } P_{W2}}$

P_{W1} is of the same form as P_{W3} when the reed is in equilibrium. When P_{W1} is a maximum, $l = (l_g - t)$ must be replaced by $(l_g - t - 2a)$ and by $(l_g - t + 2a)$ when P_{W1} is a minimum. The same holds for P_{W2} . Thus we have

$$P_{W1}^{\max} = P_{W2}^{\max} = .0445 \ln \left(1 + \frac{1.73t}{l_g - t - 2a} \right) \quad (54)$$

$$P_{W1}^{\min} = P_{W2}^{\min} = .0445 \ln \left(1 + \frac{1.73t}{l_g - t + 2a} \right) \quad (55)$$

(16) $\underline{P_{L1} \text{ and } P_{L2}}$

The term in P_{L1} arising from the inner flux path is the same as the corresponding term in P_{L3} except that again the quantity $(l_g - t)$ is replaced by $(l_g - t - 2a)$ or $(l_g - t + 2a)$ to obtain P_{L1}^{\max} and P_{L1}^{\min} , respectively.

The term arising from the outer flux path is similarly derived from the corresponding term in P_{L3} by the same substitution but an additional alteration is necessary. The portion of the reed projecting beyond the limits of the gap is triangular rather than rectangular as is the projection at the anchored end of the reed. For this reason the flux lines emanating from the walls of the pole piece have only one half the reed area in which to terminate. To take this into account it seems reasonable to substitute in the expression for y_0 , one half of the reed width for the full reed width. Thus, in Eq. (42) the value of .270 is replaced by .135. After all these adjustments are made the expression for P_{L1} and P_{L2} becomes

$$\begin{aligned} P_{L1}^{\max} = P_{L2}^{\max} &= .215 \ln \left(1 + \frac{.30}{l_g - t - 2a} \right) \\ &+ .170 - .0916 (l_g - t - 2a) \ln \frac{.30}{l_g - t - 2a} \end{aligned} \quad (56)$$

There are other permeances associated with flux paths going from corner to corner, edge to corner and face to corner. In this problem, these are not of standard form and are very difficult to estimate. They tend however, to be small compared to most of the calculated permeances, and the labor of any attempt to estimate them would not be justified by the precision attainable by our computational techniques.

CALCULATIONS OF GAP FLUX DENSITIES IN THE ABSENCE OF A REED

We now have expressions for all of the permeances relevant to the desired power calculations. Before a sample calculation for particular values of l_g , t and a are made, it would be instructive to calculate the field in the gaps when the reed is not present. Since a measured field value exists for a gap width of .060, a comparison with the result of our calculation would give an indication as to the general efficacy of the computational methods used. To do this we need the permeance P_{G4} which is given by:

$$P_{G4} = \frac{(2)(\text{Area Pole Faces})}{\text{Gap Length}} = \frac{(2)(.0219)}{.06} = .730'' \quad (57)$$

We now also have the permeances P_{E4} between the corresponding edges of the pole faces and P_{C4} between the corresponding corners on either side of the gap. P_{E4} is given by the usual formula $P_{E4} = 2(.26)y$, where y is the perimeter of a pole face and the factor of 2 enters because there are two gaps. P_{C4} is given by $(8)(.077)l_g$, where the factor 8 occurs because there are four corner-to-corner permeances for each of the two gaps. Therefore, P_{C4} is given by

$$P_{C4} = (8)(.077)(.060) = .0369'' \quad (58)$$

$P_{E4} = (.52)(.766) = .398$. P_{P4} , P_{Y4} and P_{M4} are the same as when the reed is present. We, therefore, need only to find the new values of P_{W4} and P_{L4} . P_{W4} is given by Eq. (26) with t set equal to zero. Thus

$$P_{W4} = .134 \ln \frac{.4 + (1.16)(.06)}{1.16(.06)} = .134 \ln 6.75$$

$$P_{W4} = (.134)(1.91) = .256'' \quad (59)$$

The contributions of the two outer surfaces to P_{L4} are given by the same formula as when the reed is present, but now $l = l_g = .06$ and

$$\bar{y} = \frac{.5 + .313}{2} = .406. \text{ So these contributions}$$

$$P_{L4}^0 = 2(.318)(.406) \ln \left(1 + \frac{2(.2)}{.06} \right) = .258 \ln 7.67$$

$$P_{L4}^0 = (.258)(2.04) = .526'' \quad (60)$$

In the absence of the reed the inner surfaces also contribute. The longest semicircular arc of this path, however, must be no longer than the distance d as shown in Fig. 8, so

$$\pi \left(\frac{l}{2} + w \right) = d \quad (61)$$

$$w = \frac{d}{\pi} - \frac{l}{2} \quad (62)$$

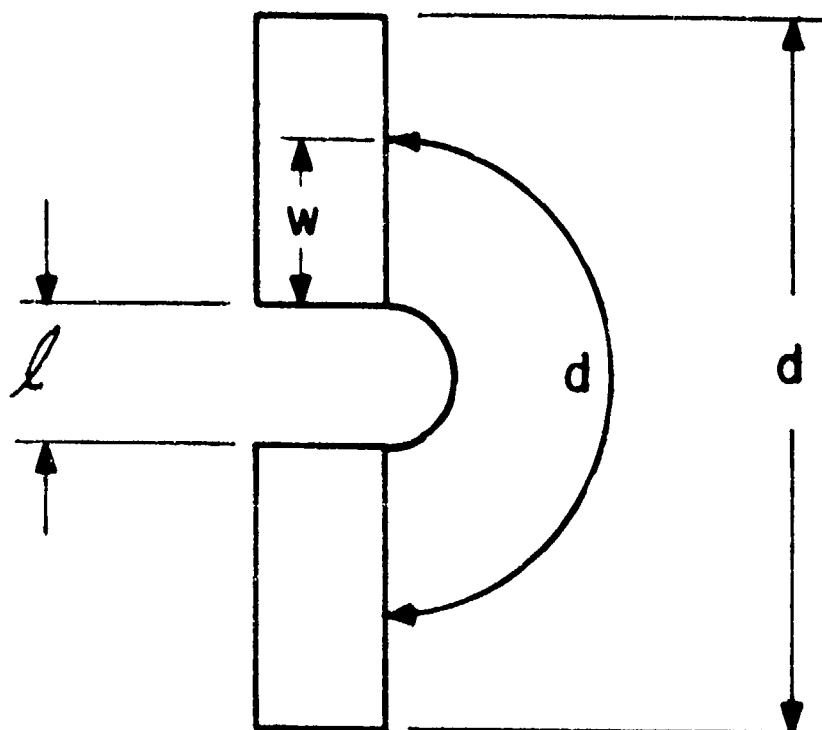


Fig. 8. Determination of the Inner Flux Paths Comprising Permeance P_{M4}

$$w = \frac{.395}{3.14} - \frac{.060}{2}$$

$$w = .126 - .030 = .0958''$$

and

$$\bar{y} = \frac{2(.313) + (2)(.0958) \tan 30^\circ}{2} = .313 + .055 = .368'' \quad (63)$$

Thus

$$P_{L4}^1 = 2(.318)(.368) \ln \left(1 + \frac{2(.0958)}{.06} \right) = .234 \ln (4.03) = .325 \quad (64)$$

$$P_{L4} = P_{L4}^1 + P_{L4}^0 = .325 + .526 = .851''$$

Our total permeance P_4 in the absence of a reed is then

$$P_4 = P_{G4} + P_{M4} + P_{Y4} + P_{H4} + P_{W4} + P_{L4} + P_{E4} + P_{C4} \quad (65)$$

$$P_4 = .730 + .647 + .212 + .731 + .256 + .851 + .398 + .037$$

$$P_4 = 3.86 \quad (66)$$

In this case P_4 is our total permeance P_t and from it we can obtain our load line slope

$$\frac{B}{H} = \frac{L}{A} P_t = \frac{(.395)(3.86)}{(.337)} = 4.52 \quad (67)$$

where $L = .395$ in and $A = .337$ in². And from the intersection of the Alnico 5 demagnetization curve with the load line $B/H = 4.52$ in Fig. 9 we obtain

$$\bar{B}_{mag} = 2670 \text{ gauss} \quad (68)$$

Since the magnetic configuration is of irregular shape, B is not perfectly uniform over the entire magnet, therefore, \bar{B}_{mag} represents an effective average flux density. So the total flux output of the magnet is

$$\Phi_t = \bar{B}_{mag} A_{mag} = (2670)(2.17) = 5790 \text{ Maxwells} \quad (69)$$

The flux in one gap Φ_G is given by

$$\Phi_G = \frac{P_{G4}}{2P_t} \Phi_t = \frac{.730}{2(3.86)} 5790 = 547 \text{ Maxwells} \quad (70)$$

and

$$B_G = \frac{\Phi_G}{A_G} = \frac{547}{.141} = 3880 \text{ gauss} \quad (71)$$

This value is about 29% higher than the measured value of 3000 gauss which is satisfactory agreement considering the various approximations made in the course of the calculations. One must also keep in mind that such comparisons should not be taken too seriously. The calculated B_G represents an average flux density obtained by dividing the total flux emanating from a pole face by the area of that pole face. The measured B_G can represent either a much more localized value of B_G or an average that includes fluxes not even associated with the gap, depending on the type, positioning and size of the probe used. The wider the gap compared to the pole face dimensions, the more difficult it is to obtain a meaningful measured value of B_G . Considering these limitations, we conclude that the general computational procedure yields reasonable agreement with the measured results, and that its judicious application to the determination of expected power outputs should not lead us too far astray.

SAMPLE CALCULATION FOR POWER OUTPUT IN THE LOAD RESISTANCE R_L

We will now perform a sample calculation for the most efficient configuration measured thus far. The relevant parameters for this arrangement are: $l_g = .070''$, $t = .032''$, $a = .017''$.

We will first calculate P_4 . We need only the components, P_{w4} and P_{L4} . Using Eq. (26) we obtain

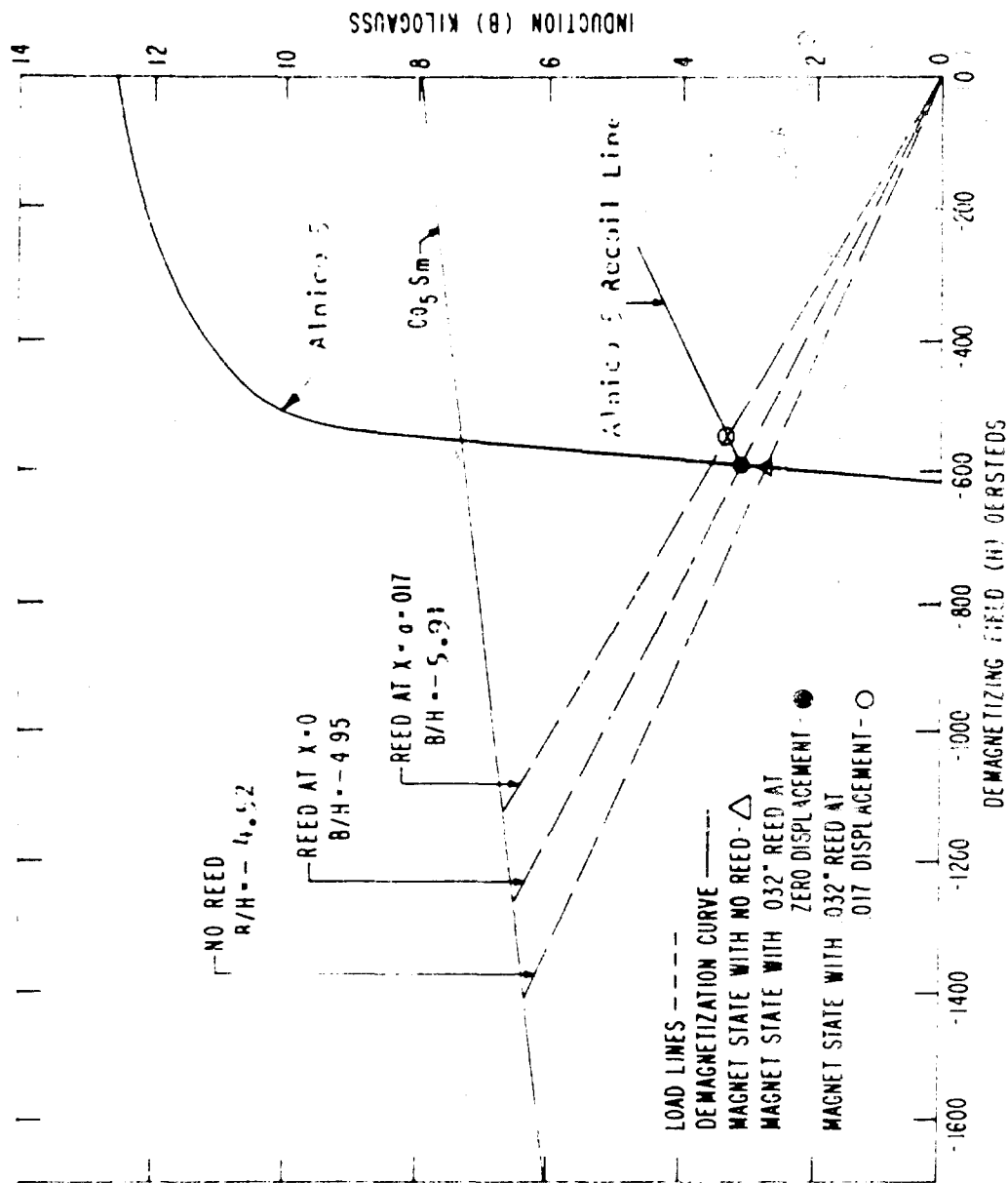


FIG. 9 DETERMINATION OF B FOR A FLUIDIC GENERATOR WITH A .070 " GAP LENGTH.

$$P_{w4} = .134 \ln \frac{.4 + 1.16 \ell_g}{t + 1.16 \ell_g} = .134 \ln \frac{.4 + (1.16)(.07)}{.032 + (1.16)(.07)}$$

$$P_{w4} = .134 \ln 4.26 = .134(1.45) = .194'' \quad (72)$$

The appropriate P_{L4} has already been calculated on page 11 and found to be .136. The other components of P_4 are the same as were used to calculate the gap fields except that now there is no P_{G4} , P_{E4} or P_{C4} . Thus

$$P_4 = P_{p4} + P_{y4} + P_{M4} + P_{L4} + P_{w4} \quad (73)$$

$$P_4 = .647 + .212 + .731 + .136 + .194$$

$$P_4 = 1.92'' \quad (74)$$

P_3 has five components

$$P_{G3} = \frac{.0408}{.070 - .032} = \frac{.0408}{.038} = 1.07'' \quad (75)$$

$$P_{E3} = .0364'' \quad (76)$$

$$P_{F3} = .303'' \quad (77)$$

$$P_{L3} = .215 \ln \left(1 + \frac{.30}{\ell_g - t} \right) + (.213 - .0916 (\ell_g - t)) \ln \left(\frac{.30}{(\ell_g - t)} \right) \quad (78)$$

$$P_{L3} = .215 \ln \left(1 + \frac{.30}{.038} \right) + (.213 - .0916(.038)) \ln \left(\frac{.30}{.038} \right)$$

$$P_{L3} = .215 \ln 8.90 + (.213 - .0035) \ln 7.90 = (.215)(2.19) + (.209)(2.07)$$

$$P_{L3} = .470 + .432''$$

$$P_{L3} = .902'' \quad (79)$$

$$P_{w3} = .0445 \ln \left(1 + \frac{1.73t}{\ell_g - t} \right) = .0445 \ln \left(1 + \frac{(1.73)(.032)}{.038} \right) \quad (80)$$

$$P_{w3} = .0445 \ln (2.46) = (.0445)(.900) = .0400'' \quad (81)$$

$$P_3 = 1.07 + .0364 + .303 + .902 + .0400$$

$$\underline{P_3 = 2.35''} \quad (82)$$

P_1 and P_2 , when the reed is at its equilibrium position, are obtained by substituting $a = 0$ in the expression for the components of P_1 and P_2 and then adding these components. This procedure yields $P_1 = P_2 = 2.27''$. We need to find P_t^E , that is, the total permeance at equilibrium to locate the point on the Alnico 5 demagnetization curve at which the recoil line begins. (Point E in Fig. 9)

$$P_t^E = P_4 + \frac{P_3}{2} + \frac{P_1}{2} = 4.23 \quad (83)$$

$$\left(\frac{B}{H}\right)_E = -\frac{l}{A} P_t^E = - (1.17)(4.23) = - 4.95 \quad (84)$$

$$B_H = 2950 \text{ gauss (from Fig. 9)} \quad (85)$$

$$H = -\frac{2950}{4.95} = - 596 \text{ oe} \quad (86)$$

The slope of recoil lines for Alnico 5 is 4.3 so the equation of the recoil line is

$$B = 4.3H + B_0 \quad (87)$$

$$2950 = (4.3)(- 596) + B_0$$

$$2950 + (4.3)(596) = B_0$$

$$2950 + 2560 = B_0$$

$$5510 = B_0 \quad (88)$$

So we have for our recoil line

$$B = 4.3H + 5510 \text{ gauss} \quad (89)$$

We now need only P_1 and P_2 to complete our analysis. Using Eqs. (50) and (51) we have

$$P_{G1}^{\max} = P_{G2}^{\max} = \frac{.0408}{l_g - t - 2a} = \frac{.0408}{.070 - .032 - 2(.017)} = \frac{.0408}{.004} \quad (90)$$

$$P_{G1}^{\max} = P_{G2}^{\max} = 10.2'' \quad (91)$$

$$P_{G2}^{\min} = P_{G1}^{\min} = \frac{.0408}{l_g - t + 2a} = \frac{.0408}{.070 - .032 + .034} = \frac{.0408}{.072} = .566'' \quad (92)$$

From (33) and (35) respectively, we have

$$P_{E1} = P_{E2} = .0364'' \quad (93)$$

$$P_{F1} = P_{F2} = .303'' \quad (94)$$

From (54)

$$\left. \begin{array}{l} p_{w1}^{\max} \\ p_{w2}^{\min} \end{array} \right| = .0445 \ln \left(1 + \frac{1.73t}{\frac{L}{g} - t \mp 2a} \right) = .0445 \ln \left(1 + \frac{(1.73)(.032)}{.070 - .032 \mp 2(.017)} \right)$$

$$p_{w1}^{\max} = .0445 \ln \left(1 + \frac{0.0554}{.004} \right) = (.0445) \ln 14.85$$

$$p_{w1}^{\max} = (.0445) (2.70) = .120' \quad (95)$$

$$p_{w2}^{\min} = .0445 \ln \left(1 + \frac{0.0554}{.072} \right) = .0445 \ln 1.77$$

$$p_{w2}^{\min} = (.0445)(.571) = .0254 \quad (96)$$

From (56)

$$p_{L1}^{\max} = .215 \ln \left(1 + \frac{.30}{.004} \right) + [.170 - (.0916)(.004)] \ln \left(\frac{.30}{.004} \right)$$

$$p_{L1}^{\max} = .215 \ln 76 + .170 \ln 75 = (.215)(4.33) + (.170)(4.32)$$

$$p_{L1}^{\max} = .931 + .734 = 1.66 \quad (97)$$

Similarly

$$p_{L2}^{\min} = .215 \ln \left(1 + \frac{.30}{.072} \right) + [.170 - (.0916)(.072)] \ln \left(\frac{.30}{.072} \right)$$

$$p_{L2}^{\min} = .215 \ln 5.17 + .163 \ln 4.17 = .215(1.64) + (.163)(1.43)$$

$$p_{L2}^{\min} = .353 + .233 = .586'' \quad (98)$$

Finally, we have for p_1^{\max} and p_2^{\min}

$$p_1^{\max} = 10.2 + .036 + .303 + .120 + 1.663$$

$$p_1^{\max} = 12.32' \quad (99)$$

$$p_2^{\min} = .566 + .036 + .303 + .025 + .586$$

$$p_2^{\min} = 1.52'' \quad (100)$$

$$p_t^{\max} = p_4 + \frac{(p_1^{\max} + p_3)(p_2^{\min} + p_3)}{p_1^{\max} + p_2^{\min} + 2p_3} \quad (101)$$

$$p_t^{\max} = 1.92 + \frac{(12.32 + 2.35)(1.52 + 2.35)}{12.32 + 2.35 + 1.52 + 2.35}$$

$$p_t^{\max} = 1.92 + \frac{(14.67)(3.87)}{18.5} = 4.98'' \quad (102)$$

$$\left(\frac{B}{H}\right)^{\max} = -\frac{i}{A} p_t^{\max} = - (1.17)(4.98) = - 5.83 \quad (103)$$

From (89)

$$B^{\max} = 4.3H + 5510$$

$$B^{\max} = \frac{5510}{1 + \frac{4.3}{5.83}} = \frac{5510}{1.73} = 3170 \text{ gauss} \quad (104)$$

$$\Phi_t^{\max} = B^{\max} A_{\text{magnet}} = (3170)(2.18) = 6880 \text{ Maxwells} \quad (105)$$

Substituting all of our calculated permeance values and $\Phi_t^{\max} = 6880$ in Eq. (11) yields for the amount of flux carried by the reed

$$\Phi_R^{\max} = 6880 \frac{2.35(12.32 - 1.52)}{(14.7)(3.87) + (1.92)(18.5)}$$

$$\Phi_R^{\max} = \frac{(6880)(2.35)(10.8)}{92.4} = 1890 \text{ Maxwells} \quad (106)$$

The induced voltage amplitude is then given by

$$V^{\max} = \frac{N L \Phi_R^{\max}}{10^8} \quad (107)$$

where ω is the angular frequency of vibration of the reed, which in this case is $2\pi(1500)$ and $N = (1500)$ is the number of turns in the coil.

$$V^{\max} = \frac{(1500)(2\pi)(1500)(1890)}{10^8} = 267 \text{ volts} \quad (108)$$

Attached to the coil is a 600- Ω load resistor and the resistance of the coil is 600 Ω as well. The inductance of the coil with a .017" reed is about 23 mhenrys. To find the power dissipated in the load, we need the total impedance Z of the circuit.

$$Z^2 = R^2 + X_L^2 = (1200)^2 + X_L^2 \quad (109)$$

The inductive reactance X_L is just

$$X_L = 2\pi fL = \frac{(6.28)(1500)(23)}{1000} = 217 \Omega \quad (110)$$

$$Z^2 = (1200)^2 + (217)^2 = 1.49 \times 10^6 \quad (111)$$

$$Z = 1220 \Omega \quad (112)$$

The root mean square current \bar{I} is given by $\bar{I} = \frac{I_{\max}}{\sqrt{2}} = \frac{V_{\max}}{Z/\sqrt{2}}$ and the power dissipated in the load p_L is then

$$p_L = \bar{I}^2 R_L = \frac{V_{\max}^2 R_L}{2Z^2} = \frac{(267)(267)(600)}{(2)(1220)(1220)} = 14.4 \text{ watts} \quad (113)$$

The actual power measured, however, was only 5 watts. The foregoing calculations were made on the assumption that the reed had infinite permeance. The electrical steel of which the reed is composed saturates at 19,900 gauss and its cross-section area is

$$A = (.27 \times .032 \times 2.54 \times 2.54) \text{ cm}^2 \quad (114)$$

Therefore, the maximum amount of flux it can carry is

$$\Phi_s = (19,900)(.27)(.032)(2.54)(2.54) = 1110 \text{ Maxwells} \quad (115)$$

and if, as was shown, 1890 Maxwells in the reed would yield 14.4 watts then the maximum allowable flux of 1110 Maxwells would give us

$$p_s = \left(\frac{1110}{1890} \right)^2 14.4 = 4.97 \text{ watts which is almost exactly the power observed.}$$

It, therefore, appears that to utilize the full flux output of the magnets used, the reed thickness must be increased by about 70%. Whether this can be done without drastically affecting the vibrational response to the oscillator is a mechanical problem beyond the scope of this study.³ It is clear, however, that increased output with the present configuration depends largely upon increasing the flux carrying capability of the reed. If, on the other

3. Carl J. Campagnuolo and Richard N. Gottron, "The Fluidic Generator: A New Electrical Power Source," 24th Annual Proceedings Power Sources Symposium, May 19-21, 1970.

hand, the present output of 5 watts is satisfactory, it can be attained with considerably less magnetic material since with the present arrangement 70% more flux is available than can be profitably used. Of course, if the reed thickness is increased, the gaps must be lengthened by an appropriate amount to keep the load line of the magnets the same and to maintain reasonable space in which the reed can oscillate. Also the expected power output would be altered slightly since the inductance of the coil is a function of reed thickness. This effect, however, would not be significant.

Parts of the keeper may also be saturated under the given operating conditions. The cross sections which would saturate first are the ones indicated in Fig. 3. Under the given conditions, the flux that must be carried by the narrowest sections of the keeper is given by

$$\Phi_s = \frac{\Phi_t^{\max}}{2} \frac{p_t^{\max} - p_m^4 - p_y^4}{p_t^{\max}} \quad (116)$$

$$\Phi_s = \frac{6880}{2} \frac{4.98 - .731 - .212}{4.98}$$

$$\Phi_s = \frac{6880}{2} \frac{4.04}{4.98} = 2790 \text{ Maxwells} \quad (117)$$

and

$$B = \frac{\Phi}{A} = \frac{2790}{(.36)(.07)(2.54)(2.54)} = 17,200 \text{ gauss} \quad (118)$$

Since permalloy saturates at 16,000 gauss, the keeper cross section need be increased by only 7% to take full advantage of the available flux at the given reed displacement. This change alone, of course, would do no good unless the reed thickness were doubled as well, as we have already shown.

EXPERIMENTAL CHECK OF CALCULATIONS

To test the validity of the foregoing calculations, a series of measurements was made to determine the maximum voltage output as a function of reed displacement under conditions for which there is no magnetic saturation. The reed displacement was measured by "stopping" the vibrating reed at its maximum displacement (a) with a strobe light and then measuring (a) with a telemicroscope equipped with a calibrated micrometer screw.

The results of this experiment are shown in Fig. 10 where they are compared with the theoretical curve plotted from results of the computer calculations listed in Appendix II. As is apparent from Fig. 10, agreement is excellent considering the many approximations used in the calculations and that (a) can be measured with a precision of no better than ten to fifteen percent.

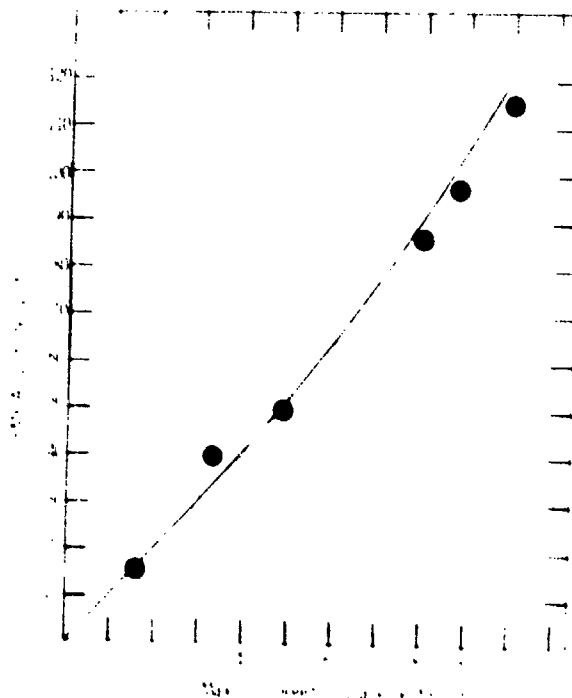


Fig. 10. Comparison of Experimental Points (dots) with Theoretical EMF-Displacement Curve (solid curve)

LOSSES IN THE REED

It is of interest to know how much electrical energy is dissipated in the reed through eddy current losses and hysteresis. The expression for the eddy current power dissipation in watts (see Appendix I) is

$$\frac{P_E}{n} = \frac{n^2 f^2 B_m^2 \tau^3 / w}{60 \cdot 10^{16}} = \frac{2 f^2 B_m^2 \tau^2 V}{60 \cdot 10^{16}} \quad (119)$$

where n is number of laminations, f is the frequency, B_m is the flux density amplitude in gauss, τ is the lamination thickness, l and w are the reed length and width respectively, V is the lamination volume, and ρ is the resistivity in ohm-cm. The reed dimensions are all measured in centimeters and the power is in watts. Since the reed is composed of two laminations each of the thickness $\tau = .0406$ cm, we multiply (119) by two and insert .0406 for τ as well as the appropriate values of the other parameters under saturation conditions.

$$P_E = .34 \text{ watts} \quad (120)$$

The expression for hysteresis power loss in the reed is given by

$$P_H = \pi f^2 \left(\frac{B_m}{10^8} \right)^{1.6} \Delta wt = \pi f^2 \left(\frac{B_m}{10^8} \right)^{1.6} V \quad (121)$$

where $V = .0706 \text{ cc.}$ is a constant depending on the material and is .002 for our reed. Making the proper insertions gives us

$$P_H = (.002)(1500)(1500) \left(\frac{19.9}{10^5} \right)^{1.6} (.0706)$$

$$P_H = \frac{(.002)(1500)(1500)(1200)(.0706)}{10^{10}}$$

$$P_H = .00004 \text{ watts which is negligible.} \quad (122)$$

The eddy current losses, however, could become quite significant if the reed were not properly laminated. For example, an unlaminated reed of the same thickness as the laminated one would have four times the eddy current loss. It is, therefore, important that if the reed thickness is increased as recommended, it be done by adding laminations. A solid unlaminated reed of twice the thickness of the one presently used would have a power loss of about .11 watts, whereas a fourfold lamination of such a reed would reduce the loss to about .68 watts.

LOSSES IN THE KEEPER

Eddy current losses in the keeper come about because of fluctuations in the flux density as the magnet load line changes with the vibration of the reed. If the keeper thickness is increased by the required 7%, or better still, widened at the narrow points which saturate by 7%, the maximum flux density would be no more than 16,000 gauss. The operating conditions would then be optimized and it would be useful to have some estimate as to the eddy current losses. Such an estimate would be very rough since it is very difficult to determine the flux pattern in the keeper. We already know that when the reed is at its greatest displacement the flux carried by the narrowest portions of the keeper would be 16,000 gauss. We also need the flux density when the reed is in its equilibrium position. This is obtained by substituting the appropriate equilibrium values of the parameters in Eqs. (116 - 118). In this fashion we obtain

$$\phi_K^E = \frac{\phi_t^E}{2} \frac{P_t^E - P_m^E - P_y^E}{P_t^E}$$

$$\phi_K^E = \frac{(2950)(2.17)}{2} \frac{4.23 - .731 - .212}{4.23} = 2490 \text{ Maxwells} \quad (123)$$

The cross-section dimensions of the keeper are .914 cm and .178 cm, so

$$B_K^E = \frac{2490}{(.914)(.178)} = 15,300 \text{ gauss}$$

Therefore, the field at the narrowest portion of the keeper fluctuates from 15,300 gauss to 16,000 gauss with an amplitude of $\frac{700}{2} = 350$ gauss. The fluctuation in other parts of the keeper would be somewhat smaller, but 350 gauss can be used to get an upper limit for the expected losses. To this end we use Eq. (119) with the appropriate constants for the keeper whose volume is 2.44 cc and whose thickness would be .178 cm. Further, the fluctuation frequency of B_K^E is double that of the reed since B_K^E is a maximum at both ends of the reed's vibrational cycle. Using these values in Eq. (119) yields:

$$P_E = \frac{(3.14)(3.14)(3000)(3000)(350)(350)(.178)(.178)(2.44) 10^6}{(6)(45) 10^{16}}$$

$$P_E = .31 \text{ watts}$$

(124)

which is about 2% of the expected optimum output. This loss of electrical energy can be avoided by a double lamination of the keeper to cut the loss to 0.08 watts or even a quadruple lamination which would result in the low eddy current loss of 0.02 watts. Since these numbers represent only upper limits, the actual values would probably be much smaller. The use of high-resistivity, high-permeability composite castings for the keepers would probably eliminate the eddy current losses entirely.

The hysteresis loss in the keeper cannot be obtained directly from Eq. (121) because the field in the keeper varies only over positive values of B , while (121) applies only to complete cycles about the principal hysteresis loops. We can, however, get some idea as to the order of magnitude of the loss by substituting our amplitude 350 gauss into (121). Using the value $\eta = .0001$ for permalloy we have

$$P_H = (.0001)(3000)(3000) \left(\frac{350}{10^8} \right)^{1.6} (2.44) \quad (125)$$

$$P_H = 4.1 \times 10^{-6} \text{ watts} \quad (126)$$

which as in the case of the reed is negligible.

SUMMARY AND RECOMMENDATIONS

(1) Under the present conditions of operation the fluidic generator is producing only about 35% of its potential power output. Doubling the cross-section of the reed by adding lamination or increasing its width would augment its flux carrying capability to where the potential power output could be fully realized with only negligible eddy current and hysteresis losses in the reed. Appropriate adjustments in gap length would also be necessary to prevent saturation and to allow sufficient latitude for vibration.

(2) Keeper hysteresis losses are also negligible but eddy current losses can be appreciable unless the keeper is laminated or made of a suitable composite material. Further, the narrowest part of the keeper cross section must be increased by about 7%, to avoid saturation with the given reed displacement amplitude of .017".

(3) To facilitate design optimization of future generators a computer study was done to obtain a series of design matrices which list voltage amplitude as a function of gap length (l_g), reed thickness (τ), and reed displacement amplitude (a). The computer program uses the mathematical procedure of the sample calculation made in this report and it is written in Fortran IV for the Burroughs 5500 Computer. The program is reproduced in Appendix II together with a summary of the equations used in the calculations.

The computed voltage amplitudes are listed in Table II in Appendix II. All of the listed voltages are calculated for a 1500 turn coil, a reed width of .27", a reed vibrational frequency of 1500 hertz, and the standard magnet and keeper dimensions used in the sample calculation and summarized in Table I. The assumption is also made that the keeper does not saturate for any of the parameter choices listed. This assumption, as we have shown in the sample, is not valid for the larger values of (a) unless the keeper thickness is increased by an appropriate amount. This thickening of the keeper would, of course, affect the value of P_{y4} and hence the permeance calculations. The effect on the calculated voltage output, however, is not likely to be significant because in the sample calculation P_{y4} is only of the order of 10% of P_4 , and P_4 in turn constitutes only about 40% of P_t . Furthermore, P_{y4} itself is relatively insensitive to keeper thickness (W) since W appears only in a small correction to the logarithmic factor in Eq. (13). Hence, thickening of the keeper to required values will probably not affect the calculated voltages by more than the order of a percent or two even in the worst cases.

The calculations also assume infinite reed permeability and the values not actually attainable with the reed material used are marked with an asterisk. Thus the design can easily be optimized by choosing the parameters yielding the largest non-asterisked voltage in the Tables. This, of course, is always subject to the condition that the reed of requisite thickness can be driven with the stated amplitude and frequency. Perhaps this can be assured in many cases by making appropriate adjustments of the reed stiffness through variation of the reed material and dimensions at the point where it is anchored or by widening the reed rather than thickening it.

The Tables can be used to obtain the power output in any load resistance R_L from Eq. (126), viz.

$$P_L = \frac{V_{\max}^2 R_L}{2Z^2} \quad (127)$$

where Z is the total circuit impedance obtained from Eqs. (109) and (110). The coil inductance is a function of reed thickness and the empirical curve of the relationship is shown in Fig. 11. As is evident from (111), however, the reactive portion of the impedance is only about a percent and a half of the total, and the variation of reactance with reed thickness will not have a significant effect upon power output at the frequency and resistance used.

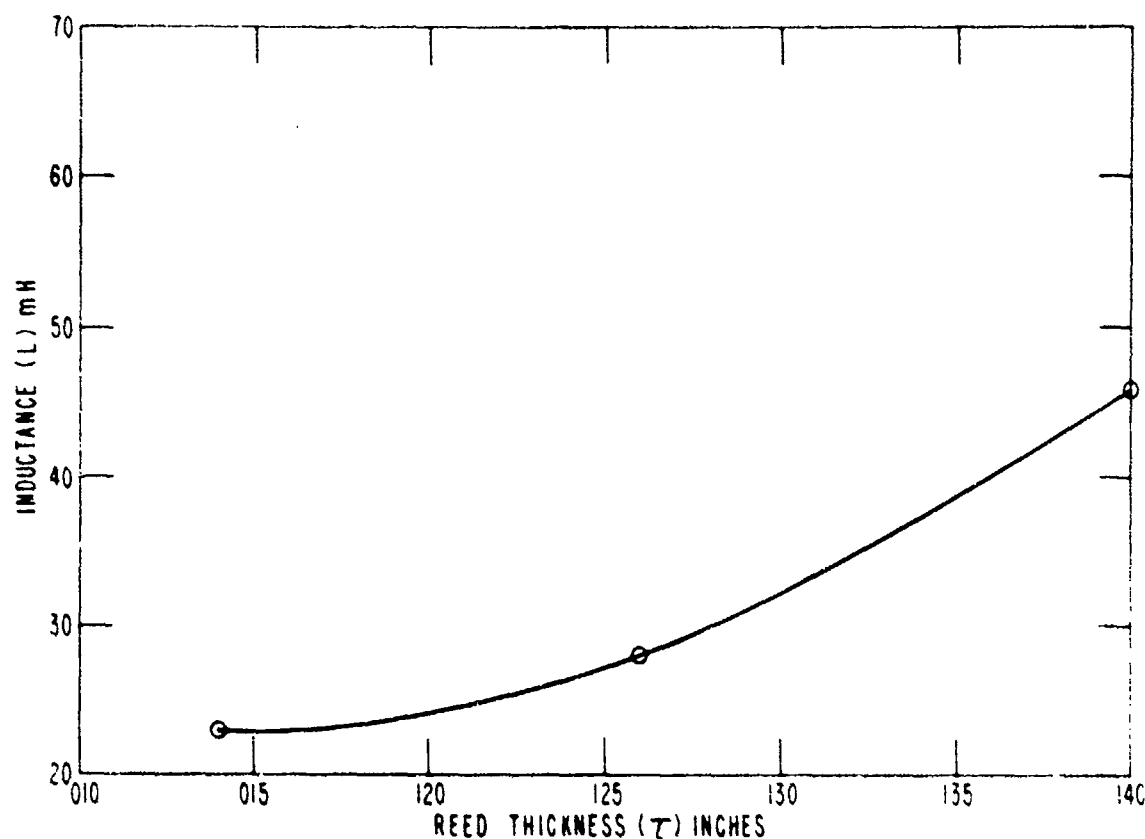


Fig. 11. Inductance of Sensing Coil as a Function of Reed Thickness.

(4) The present configuration of the magnetic circuit results in relatively low-sloped load lines which intersect the Alnico 5 demagnetization curve well below its 'knee' (Fig. 9). Note that if Co_5Sm were used in place of Alnico 5 much higher flux densities would result for any given load line with a concomitant increase in power output. This, of course, is always assuming that neither reed nor keeper are saturated. It probably would be most practical to keep the power output at about 10-12 watts with the use of much less magnetic material through substitution of Co_5Sm . A more compact generator operating at lower load lines with the same power output could probably be designed using the latter material.

(5) A composite material consisting of permalloy particles densely packed in a nonconducting binder might be desirable for the keeper. Providing a sufficiently high packing density is obtainable, eddy current losses could be virtually eliminated without too great a loss in flux-carrying capability. Such a composite could also be press molded into the desired shape thus improving the efficiency of fabrication.

APPENDIX I

DERIVATION OF THE FORMULA FOR ENERGY LOSS THROUGH EDDY CURRENTS

Consider a sheet of metal such as the one pictured in Fig. 12. The length l is much larger than the thickness τ and the width w is arbitrary. A sinusoidal field B is acting in the y direction as shown. The varying field B will induce EMFs which will give rise to circulating or eddy currents. We wish to find the power expended by these currents. We proceed by considering a current loop indicated by the arrows in Fig. 12. The EMF around this loop is then given by

$$V = \frac{d\Phi}{10^8 dt} = \frac{A dB}{10^8 dt} = \frac{2lx}{10^8} \frac{dB}{dt} \quad (128)$$

If we neglect the small voltage drops across the width of the loop at the ends we can write for the voltage drop per unit length in the z direction

$$v = \frac{V}{2l} = \frac{x dB}{10^8 dt} \quad (129)$$

The current density i is then obtained from the relation

$$i = \frac{v}{\rho} \quad (130)$$

where ρ is the resistivity of the material. Substituting (129) in (130) yields

$$i = \frac{x}{10^8 \rho} \frac{dB}{dt} = \frac{x B_{\max}}{\rho 10^8} \cos \omega t \quad (131)$$

The instantaneous power, dP , dissipated in the unit laminar element dx is then

$$dP = w i^2 dx = \frac{w x^2 \omega^2 B_{\max}^2}{\rho 10^{16}} \cos^2 \omega t$$

$$dP = \frac{w \omega^2 B_{\max}^2}{10^{16} \rho} x^2 dx \cos^2 \omega t \quad (132)$$

and the total power loss, P , of a unit length of sheet is

$$P = \int_{-\frac{\tau}{2}}^{\frac{\tau}{2}} dP = \frac{w \omega^2 B_{\max}^2}{3 \rho 10^{16}} x^3 \bigg|_{-\frac{\tau}{2}}^{\frac{\tau}{2}} \cos^2 \omega t \quad (133)$$

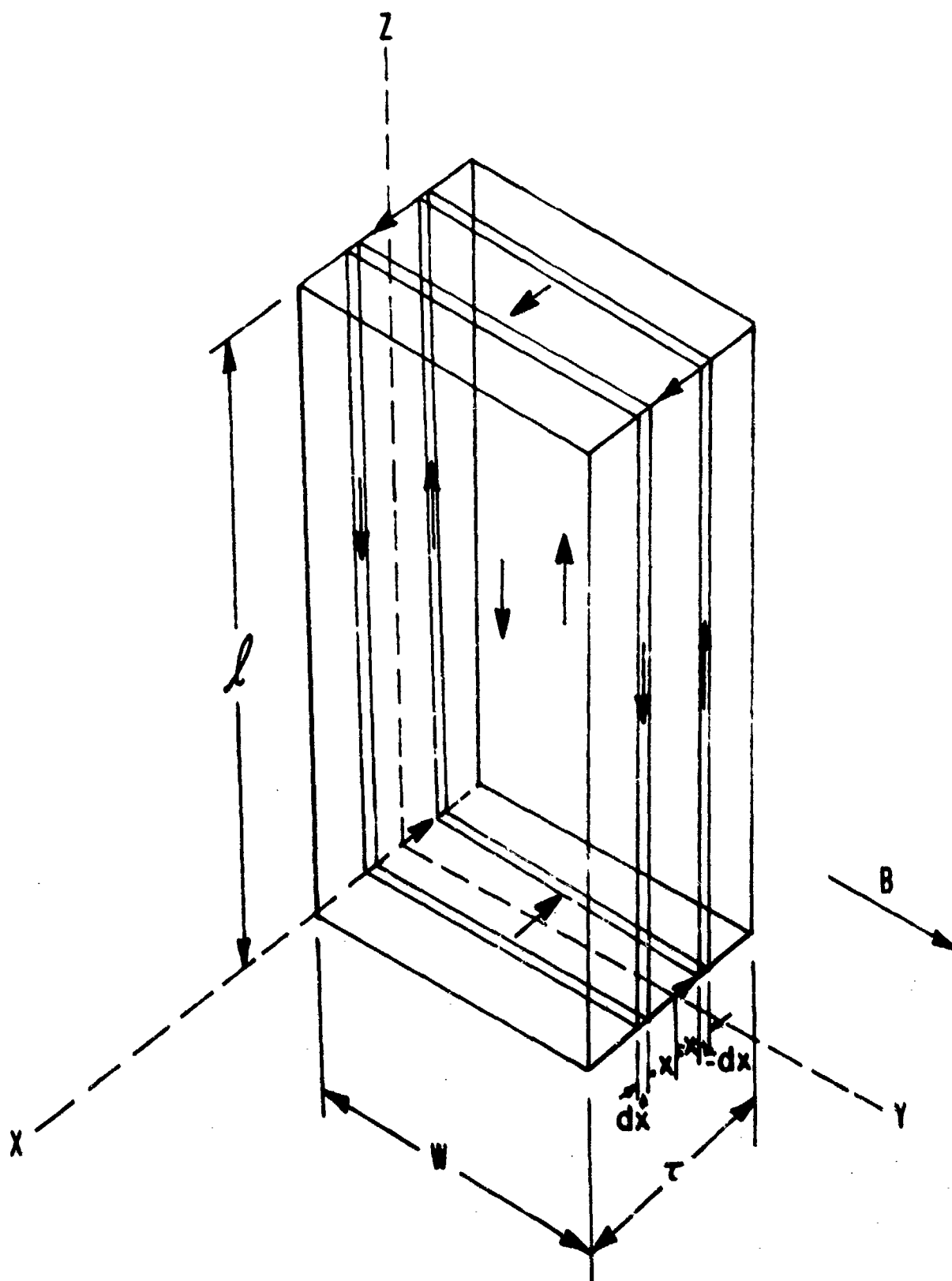


Fig. 12. Derivation of the Eddy Current Loss Formula.

$$P = \frac{w w_B^2 \tau^3}{3 \rho 10^{16} 4} \cos^2 \omega t = \frac{w w_B^2 \tau^3}{12 \rho 10^{16}} \cos^2 \omega t \quad (134)$$

The energy loss per cycle is then given by

$$W_c = \frac{w w_B^2 \tau^3}{12 \rho 10^{16}} \int_0^{2\pi} \cos^2(\omega t) (\omega dt) \quad (135)$$

$$W_c = \frac{\pi w w_B^2 \tau^3}{12 \rho 10^{16}} \quad (136)$$

The average power per unit length of sheet is

$$\bar{P} = f W_c = \frac{f \pi w w_B^2 \tau^3}{12 \rho 10^{16}} = \frac{\pi^2 w f^2 \tau^3 w_B^2}{6 \rho 10^{16}} \quad (137)$$

The average power for the entire length of the sheet is then

$$\bar{P}_t = L \bar{P} = \frac{\pi^2 f^2 \tau^2 w_B^2 L w \tau}{6 \rho 10^{16}} = \frac{\pi^2 f^2 \tau^2 w_B^2 V}{6 \rho 10^{16}} \quad (138)$$

where V is the volume of the sheet.

APPENDIX II

SUMMARY OF COMPUTATIONAL PROCEDURE IN CALCULATION OF PEAK VOLTAGE AND ADAPTATION TO COMPUTER PROGRAM

All relevant permeances are calculated for the generator with the reed in the equilibrium position. There are five such permeances of which four are associated with the four air gaps and the fifth with the magnetic leakage paths. The two gaps at the fixed end of the reed are always equal and the associated permeances are equal and denoted by P_3 . P_1 and P_2 refer to the gaps at the vibrating end of the reed. P_4 is the permeance of the composite of all leakage paths. P_3 has five components and we can write for the equilibrium condition:

$$P_3 = P_{G3} + P_{E3} + P_{F3} + P_{L3} + P_{W3}$$

all of which are illustrated in Fig. 5 and are computed by the following formulae as functions of l_g , t and a . All other dimensions are fixed and are the same as those of the sample calculations.

$$P_{G3} = \frac{.0408}{l_g - t}$$

$$P_{E3} = .0364$$

$$P_{F3} = .303$$

$$P_{L3} = .215 \ln \left(1 + \frac{.30}{l_g - t} \right) + [.213 - .0916(l_g - t)] \ln \left(\frac{.30}{l_g - t} \right)$$

$$P_{W3} = .0445 \ln \left(1 + \frac{1.73 t}{l_g} \right)$$

P_4 has five components and is given by

$$P_4 = P_{p4} + P_{y4} + P_{M4} + P_{L4} + P_{W4}$$

and

$$P_{p4} = .647''$$

$$P_{y4} = .212''$$

$$P_{M4} = .731''$$

$$P_{L4} = .136''$$

$$P_{W4} = .134 \ln \frac{.4 + 1.16 l}{t + 1.16 l_g}$$

P_1 and P_2 have the same form with either always reaching its maximum when the other is at its minimum. Each has 5 components and

$$P_1 = P_{G1} + P_{E1} + P_{F1} + P_{L1} + P_{W1}$$

$$P_{G1 \min}^{\max} = P_{G2 \min}^{\max} = \frac{.0408}{l_g - t \mp 2a}$$

$$P_{E1} = P_{E2} = .0364''$$

$$P_{F1} = P_{F2} = .303''$$

$$P_{L1 \min}^{\max} = P_{L2 \min}^{\max} = .215 \ln \left(1 + \frac{.30}{l_g - t \mp 2a} \right)$$

$$+ [.170 - .0918(l_g - t \mp 2a)] \frac{.30}{l_g - t \mp 2a}$$

$$P_{W1 \min}^{\max} = P_{W2 \min}^{\max} = .0445 \ln \left(1 + \frac{1.73 t}{l_g - t \mp 2a} \right)$$

The total permeance of the magnetic circuit P_t is then given by the formula:

$$P_t = P_4 + \frac{(P_1 + P_3)(P_2 + P_3)}{P_1 + P_2 + 2P_3}$$

and both the equilibrium and full reed displacement values of P_t are calculated. We use the equilibrium value P_t^E to obtain our load line slope B/H via the formula

$$\left(\frac{B}{H} \right)^E = - \frac{l}{A} P_t^E$$

where l and A are respectively the length and cross-sectional area of the magnet.

The intersection of this load line with the demagnetization curve of Alnico 5 is the base point of the recoil line (slope 4.3 for Alnico 5) along which the system then operates. The recoil line is thus of the form:

$$B = 4.3 H + B_0$$

where B_0 is determined from the base point. The maximum value of P_t corresponding to full displacement is then used in the relation

$$\left(\frac{B}{H} \right)^{\max} = - \frac{l}{A} P_t^{\max}$$

to obtain the new load line. The intersection of this load line with the recoil line already determined yields the maximum value of B in the magnet.

The corresponding total available flux is then given by

$$\Phi_t^{\max} = B^{\max} A_{\text{magnet}}$$

and the flux amplitude in the reed can then be calculated from the relation

$$\Phi_R^{\max} = \Phi_t^{\max} \frac{P_3(P_1 - P_2)}{(P_1 + P_3)(P_2 + P_3) + P_4(P_1 + 2P_3 + P_2)}$$

The induced voltage amplitude is then obtained from

$$V^{\max} = \frac{N\omega \Phi_R^{\max}}{10^8}$$

The computer program which follows this procedure is reproduced below in Fortran IV. In addition to calculating the expected voltage for a given set of parameters, the program also instructs the computer to calculate

$$B_R^{\max} = \frac{\Phi_R^{\max}}{10^8}$$

If the result is larger than the saturation flux density of the reed material (19,900 gauss), the calculated voltage is not attainable and is marked with an asterisk in Table I.

COMPUTER PROGRAM USED TO OBTAIN TABLE II

000000100
000000200
000000300
000000400
000000500
000000600
000000700
000000800
000000900
000001000
000001100
000001200
000001300
000001400
000001500
000001600
000001700
000001800
000001900
000002000
000002100
000002200
000002300
000002400
000002500
000002600
000002700
000002800
000002900
000003000
000003100
000003200
000003300
000003400
000003500
000003600
000003700
000003800
000003900
000004000
000004100
000004200
000004300
000004400
000004500
000004600
000004700
000004800
000004900
000005000
000005100
000005200
000005300
000005400
000005500
000005600
000005700
000005800
000005900
000006000
000006100
000006200

00006300
00006400
00006500
00006600
00006700
00006800
00006900
00007000
00007100
00007200
00007300
00007400
00007500
00007600
00007700
00007800
00007900
00008000
00008100
00008200
00008300
00008400
00008500
00008600
00008700
00008800
00008900
00009000
00009100
00009200

TABLE II

Voltage outputs for various gap lengths, reed thicknesses and reed displacement amplitudes. The values marked with an asterisk are not attainable with the reed material presently in use. They are the values to be expected with the use of reed material that does not saturate under the given operating conditions.

GENERATED OUTPUT IN VOLTS AC FOR GAP LENGTH = 0.040 INCHES

REED THICKNESS (INCHES)	AMPLITUDE IN INCHES	AMPLITUDE IN INCHES
.010	.01	.01
.012	.02	.02
.014	.03	.03
.016	.04	.04
.018	.05	.05
.020	.06	.06
.022	.07	.07
.024	.08	.08
.026	.09	.09
.028	.10	.10
.030	.11	.11
.032	.12	.12
.034	.13	.13
.036	.14	.14
.038	.15	.15
.040	.16	.16
.042	.17	.17
.044	.18	.18
.046	.19	.19
.048	.20	.20
.050	.21	.21
.052	.22	.22
.054	.23	.23
.056	.24	.24
.058	.25	.25
.060	.26	.26
.062	.27	.27
.064	.28	.28
.066	.29	.29
.068	.30	.30
.070	.31	.31
.072	.32	.32
.074	.33	.33
.076	.34	.34
.078	.35	.35
.080	.36	.36
.082	.37	.37
.084	.38	.38
.086	.39	.39
.088	.40	.40
.090	.41	.41
.092	.42	.42
.094	.43	.43
.096	.44	.44
.098	.45	.45
.100	.46	.46
.102	.47	.47
.104	.48	.48
.106	.49	.49
.108	.50	.50
.110	.51	.51
.112	.52	.52
.114	.53	.53
.116	.54	.54
.118	.55	.55
.120	.56	.56
.122	.57	.57
.124	.58	.58
.126	.59	.59
.128	.60	.60
.130	.61	.61
.132	.62	.62
.134	.63	.63
.136	.64	.64
.138	.65	.65
.140	.66	.66
.142	.67	.67
.144	.68	.68
.146	.69	.69
.148	.70	.70
.150	.71	.71
.152	.72	.72
.154	.73	.73
.156	.74	.74
.158	.75	.75
.160	.76	.76
.162	.77	.77
.164	.78	.78
.166	.79	.79
.168	.80	.80
.170	.81	.81
.172	.82	.82
.174	.83	.83
.176	.84	.84
.178	.85	.85
.180	.86	.86
.182	.87	.87
.184	.88	.88
.186	.89	.89
.188	.90	.90
.190	.91	.91
.192	.92	.92
.194	.93	.93
.196	.94	.94
.198	.95	.95
.200	.96	.96
.202	.97	.97
.204	.98	.98
.206	.99	.99
.208	1.00	1.00
.210	1.01	1.01
.212	1.02	1.02
.214	1.03	1.03
.216	1.04	1.04
.218	1.05	1.05
.220	1.06	1.06
.222	1.07	1.07
.224	1.08	1.08
.226	1.09	1.09
.228	1.10	1.10
.230	1.11	1.11
.232	1.12	1.12
.234	1.13	1.13
.236	1.14	1.14
.238	1.15	1.15
.240	1.16	1.16
.242	1.17	1.17
.244	1.18	1.18
.246	1.19	1.19
.248	1.20	1.20
.250	1.21	1.21
.252	1.22	1.22
.254	1.23	1.23
.256	1.24	1.24
.258	1.25	1.25
.260	1.26	1.26
.262	1.27	1.27
.264	1.28	1.28
.266	1.29	1.29
.268	1.30	1.30
.270	1.31	1.31
.272	1.32	1.32
.274	1.33	1.33
.276	1.34	1.34
.278	1.35	1.35
.280	1.36	1.36
.282	1.37	1.37
.284	1.38	1.38
.286	1.39	1.39
.288	1.40	1.40
.290	1.41	1.41
.292	1.42	1.42
.294	1.43	1.43
.296	1.44	1.44
.298	1.45	1.45
.300	1.46	1.46
.302	1.47	1.47
.304	1.48	1.48
.306	1.49	1.49
.308	1.50	1.50
.310	1.51	1.51
.312	1.52	1.52
.314	1.53	1.53
.316	1.54	1.54
.318	1.55	1.55
.320	1.56	1.56
.322	1.57	1.57
.324	1.58	1.58
.326	1.59	1.59
.328	1.60	1.60
.330	1.61	1.61
.332	1.62	1.62
.334	1.63	1.63
.336	1.64	1.64
.338	1.65	1.65
.340	1.66	1.66
.342	1.67	1.67
.344	1.68	1.68
.346	1.69	1.69
.348	1.70	1.70
.350	1.71	1.71
.352	1.72	1.72
.354	1.73	1.73
.356	1.74	1.74
.358	1.75	1.75
.360	1.76	1.76
.362	1.77	1.77
.364	1.78	1.78
.366	1.79	1.79
.368	1.80	1.80
.370	1.81	1.81
.372	1.82	1.82
.374	1.83	1.83
.376	1.84	1.84
.378	1.85	1.85
.380	1.86	1.86
.382	1.87	1.87
.384	1.88	1.88
.386	1.89	1.89
.388	1.90	1.90
.390	1.91	1.91
.392	1.92	1.92
.394	1.93	1.93
.396	1.94	1.94
.398	1.95	1.95
.400	1.96	1.96
.402	1.97	1.97
.404	1.98	1.98
.406	1.99	1.99
.408	2.00	2.00
.410	2.01	2.01
.412	2.02	2.02
.414	2.03	2.03
.416	2.04	2.04
.418	2.05	2.05
.420	2.06	2.06
.422	2.07	2.07
.424	2.08	2.08
.426	2.09	2.09
.428	2.10	2.10
.430	2.11	2.11
.432	2.12	2.12
.434	2.13	2.13
.436	2.14	2.14
.438	2.15	2.15
.440	2.16	2.16
.442	2.17	2.17
.444	2.18	2.18
.446	2.19	2.19
.448	2.20	2.20
.450	2.21	2.21
.452	2.22	2.22
.454	2.23	2.23
.456	2.24	2.24
.458	2.25	2.25
.460	2.26	2.26
.462	2.27	2.27
.464	2.28	2.28
.466	2.29	2.29
.468	2.30	2.30
.470	2.31	2.31
.472	2.32	2.32
.474	2.33	2.33
.476	2.34	2.34
.478	2.35	2.35
.480	2.36	2.36
.482	2.37	2.37
.484	2.38	2.38
.486	2.39	2.39
.488	2.40	2.40
.490	2.41	2.41
.492	2.42	2.42
.494	2.43	2.43
.496	2.44	2.44
.498	2.45	2.45
.500	2.46	2.46
.502	2.47	2.47
.504	2.48	2.48
.506	2.49	2.49
.508	2.50	2.50
.510	2.51	2.51
.512	2.52	2.52
.514	2.53	2.53
.516	2.54	2.54
.518	2.55	2.55
.520	2.56	2.56
.522	2.57	2.57
.524	2.58	2.58
.526	2.59	2.59
.528	2.60	2.60
.530	2.61	2.61
.532	2.62	2.62
.534	2.63	2.63
.536	2.64	2.64
.538	2.65	2.65
.540	2.66	2.66
.542	2.67	2.67
.544	2.68	2.68
.546	2.69	2.69
.548	2.70	2.70
.550	2.71	2.71
.552	2.72	2.72
.554	2.73	2.73
.556	2.74	2.74
.558	2.75	2.75
.560	2.76	2.76
.562	2.77	2.77
.564	2.78	2.78
.566	2.79	2.79
.568	2.80	2.80
.570	2.81	2.81
.572	2.82	2.82
.574	2.83	2.83
.576	2.84	2.84
.578	2.85	2.85
.580	2.86	2.86
.582	2.87	2.87
.584	2.88	2.88
.586	2.89	2.89
.588	2.90	2.90
.590	2.91	2.91
.592	2.92	2.92
.594	2.93	2.93
.596	2.94	2.94
.598	2.95	2.95
.600	2.96	2.96
.602	2.97	2.97
.604	2.98	2.98
.606	2.99	2.99
.608	3.00	3.00
.610	3.01	3.01
.612	3.02	3.02
.614	3.03	3.03
.616	3.04	3.04
.618	3.05	3.05
.620	3.06	3.06
.622	3.07	3.07
.624	3.08	3.08
.626	3.09	3.09
.628	3.10	3.10
.630	3.11	3.11
.632	3.12	3.12
.634	3.13	3.13
.636	3.14	3.14
.638	3.15	3.15
.640	3.16	3.16
.642	3.17	3.17
.644	3.18	3.18
.646	3.19	3.19
.648	3.20	3.20
.650	3.21	3.21
.652	3.22	3.22
.654	3.23	3.23
.656	3.24	3.24
.658	3.25	3.25
.660	3.26	3.26
.662	3.27	3.27
.664	3.28	3.28
.666	3.29	3.29
.668	3.30	3.30
.670	3.31	3.31
.672	3.32	3.32
.674	3.33	3.33
.676	3.34	3.34
.678	3.35	3.35
.680	3.36	3.36
.682	3.37	3.37
.684	3.38	3.38
.686	3.39	3.39
.688	3.40	3.40
.690	3.41	3.41
.692	3.42	3.42
.694	3.43	3.43
.696	3.44	3.44
.698	3.45	3.45
.700	3.46	3.46
.702	3.47	3.47
.704	3.48	3.48
.706	3.49	3.49
.708	3.50	3.50
.710	3.51	3.51
.712	3.52	3.52
.714	3.53	3.53
.716	3.54	3.54
.718	3.55	3.55
.720	3.56	3.56
.722	3.57	3.57
.724	3.58	3.58
.726	3.59	3.59
.728	3.60	3.60
.730	3.61	3.61
.732	3.62	3.62
.734	3.63	3.63
.736	3.64	3.64
.738	3.65	3.65
.740	3.66	3.66
.742	3.67	3.67
.744	3.68	3.68
.746	3.69	3.69
.748	3.70	3.70
.750	3.71	3.71
.752	3.72	3.72
.754	3.73	3.73
.756	3.74	3.74
.758	3.75	3.75
.760	3.76	3.76
.762	3.77	3.77
.764	3.78	3.78
.766	3.79	3.79
.768	3.80	3.80
.770	3.81	3.81
.772	3.82	3.82
.774	3.83	3.83
.776	3.84	3.84
.778	3.85	3.85
.780	3.86	3.86
.782	3.87	3.87
.784	3.88	3.88
.786	3.89	3.89
.788	3.90	3.90
.790	3.91	3.91
.792		

REMARKS: TOTAL VOLUME = 0.005 LITERS

42

SUBJECTS

43

TABLE 11 (cont'd)

SECRET
NOFORN

255931
1947

(Inches) 1.01.007.003.005.004.007.006.000.011.012.013.018.015.016.017.011.019.020.021.022.023.024.025.026.027.028.029

[illegible]

TABLE II (cont'd)

GENERATION OUTPUT IN VOLTS AC FOR 600 L.F.C.T.M. @ 0.040 INCHES

REF. NO.	1	2	3	4	5	6	7	8	9	10	11	12	13	14	15	16	17	18	19	20	21	22	23	24	25	26	27	28	29	30	31	32	33	34	35	36	37	38	39	40	41	42	43	44	45	46	47	48	49	50	51	52	53	54	55	56	57	58	59	60	61	62	63	64	65	66	67	68	69	70	71	72	73	74	75	76	77	78	79	80	81	82	83	84	85	86	87	88	89	90	91	92	93	94	95	96	97	98	99	100	101	102	103	104	105	106	107	108	109	110	111	112	113	114	115	116	117	118	119	120	121	122	123	124	125	126	127	128	129	130	131	132	133	134	135	136	137	138	139	140	141	142	143	144	145	146	147	148	149	150	151	152	153	154	155	156	157	158	159	160	161	162	163	164	165	166	167	168	169	170	171	172	173	174	175	176	177	178	179	180	181	182	183	184	185	186	187	188	189	190	191	192	193	194	195	196	197	198	199	200	201	202	203	204	205	206	207	208	209	210	211	212	213	214	215	216	217	218	219	220	221	222	223	224	225	226	227	228	229	230	231	232	233	234	235	236	237	238	239	240	241	242	243	244	245	246	247	248	249	250	251	252	253	254	255	256	257	258	259	260	261	262	263	264	265	266	267	268	269	270	271	272	273	274	275	276	277	278	279	280	281	282	283	284	285	286	287	288	289	290	291	292	293	294	295	296	297	298	299	300	301	302	303	304	305	306	307	308	309	310	311	312	313	314	315	316	317	318	319	320	321	322	323	324	325	326	327	328	329	330	331	332	333	334	335	336	337	338	339	340	341	342	343	344	345	346	347	348	349	350	351	352	353	354	355	356	357	358	359	360	361	362	363	364	365	366	367	368	369	370	371	372	373	374	375	376	377	378	379	380	381	382	383	384	385	386	387	388	389	390	391	392	393	394	395	396	397	398	399	400	401	402	403	404	405	406	407	408	409	410	411	412	413	414	415	416	417	418	419	420	421	422	423	424	425	426	427	428	429	430	431	432	433	434	435	436	437	438	439	440	441	442	443	444	445	446	447	448	449	450	451	452	453	454	455	456	457	458	459	460	461	462	463	464	465	466	467	468	469	470	471	472	473	474	475	476	477	478	479	480	481	482	483	484	485	486	487	488	489	490	491	492	493	494	495	496	497	498	499	500	501	502	503	504	505	506	507	508	509	510	511	512	513	514	515	516	517	518	519	520	521	522	523	524	525	526	527	528	529	530	531	532	533	534	535	536	537	538	539	540	541	542	543	544	545	546	547	548	549	550	551	552	553	554	555	556	557	558	559	560	561	562	563	564	565	566	567	568	569	570	571	572	573	574	575	576	577	578	579	580	581	582	583	584	585	586	587	588	589	590	591	592	593	594	595	596	597	598	599	600	601	602	603	604	605	606	607	608	609	610	611	612	613	614	615	616	617	618	619	620	621	622	623	624	625	626	627	628	629	630	631	632	633	634	635	636	637	638	639	640	641	642	643	644	645	646	647	648	649	650	651	652	653	654	655	656	657	658	659	660	661	662	663	664	665	666	667	668	669	670	671	672	673	674	675	676	677	678	679	680	681	682	683	684	685	686	687	688	689	690	691	692	693	694	695	696	697	698	699	700	701	702	703	704	705	706	707	708	709	710	711	712	713	714	715	716	717	718	719	720	721	722	723	724	725	726	727	728	729	730	731	732	733	734	735	736	737	738	739	740	741	742	743	744	745	746	747	748	749	750	751	752	753	754	755	756	757	758	759	760	761	762	763	764	765	766	767	768	769	770	771	772	773	774	775	776	777	778	779	780	781	782	783	784	785	786	787	788	789	790	791	792	793	794	795	796	797	798	799	800	801	802	803	804	805	806	807	808	809	810	811	812	813	814	815	816	817	818	819	820	821	822	823	824	825	826	827	828	829	830	831	832	833	834	835	836	837	838	839	840	841	842	843	844	845	846	847	848	849	850	851	852	853	854	855	856	857	858	859	860	861	862	863	864	865	866	867	868	869	870	871	872	873	874	875	876	877	878	879	880	881	882	883	884	885	886	887	888	889	890	891	892	893	894	895	896	897	898	899	900	901	902	903	904	905	906	907	908	909	910	911	912	913	914	915	916	917	918	919	920	921	922	923	924	925	926	927	928	929	930	931	932	933	934	935	936	937	938	939	940	941	942	943	944	945	946	947	948	949	950	951	952	953	954	955	956	957	958	959	960	961	962	963	964	965	966	967	968	969	970	971	972	973	974	975	976	977	978	979	980	981	982	983	984	985	986	987	988	989	990	991	992	993	994	995	996	997	998	999	1000
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68464870 3.754 16.171 11.4.11 0:34:33035

[illegible]

TABLE 11 (cont'd)

GENERATED INPUT IN VOLT AC FOR GAP LENGTH = 0.070 INCHES

INCHES	0.01	0.02	0.03	0.04	0.05	0.06	0.07	0.08	0.09	0.10	0.11	0.12	0.13	0.14	0.15	0.16	0.17	0.18	0.19	0.20	0.21	0.22	0.23	0.24	0.25	0.26	0.27	0.28	0.29	0.30
0.01	0.01	0.02	0.03	0.04	0.05	0.06	0.07	0.08	0.09	0.10	0.11	0.12	0.13	0.14	0.15	0.16	0.17	0.18	0.19	0.20	0.21	0.22	0.23	0.24	0.25	0.26	0.27	0.28	0.29	0.30
0.02	0.02	0.04	0.06	0.08	0.10	0.12	0.14	0.16	0.18	0.20	0.22	0.24	0.26	0.28	0.30	0.32	0.34	0.36	0.38	0.40	0.42	0.44	0.46	0.48	0.50	0.52	0.54	0.56	0.58	0.60
0.03	0.03	0.06	0.09	0.12	0.15	0.18	0.21	0.24	0.27	0.30	0.33	0.36	0.39	0.42	0.45	0.48	0.51	0.54	0.57	0.60	0.63	0.66	0.69	0.72	0.75	0.78	0.81	0.84	0.87	0.90
0.04	0.04	0.08	0.12	0.16	0.20	0.24	0.28	0.32	0.36	0.40	0.44	0.48	0.52	0.56	0.60	0.64	0.68	0.72	0.76	0.80	0.84	0.88	0.92	0.96	1.00	1.04	1.08	1.12	1.16	1.20
0.05	0.05	0.10	0.15	0.20	0.25	0.30	0.35	0.40	0.45	0.50	0.55	0.60	0.65	0.70	0.75	0.80	0.85	0.90	0.95	1.00	1.05	1.10	1.15	1.20	1.25	1.30	1.35	1.40	1.45	1.50
0.06	0.06	0.12	0.18	0.24	0.30	0.36	0.42	0.48	0.54	0.60	0.66	0.72	0.78	0.84	0.90	0.96	1.02	1.08	1.14	1.20	1.26	1.32	1.38	1.44	1.50	1.56	1.62	1.68	1.74	1.80
0.07	0.07	0.14	0.21	0.28	0.35	0.42	0.49	0.56	0.63	0.70	0.77	0.84	0.91	0.98	1.05	1.12	1.19	1.26	1.33	1.40	1.47	1.54	1.61	1.68	1.75	1.82	1.89	1.96	2.03	2.10
0.08	0.08	0.16	0.24	0.32	0.40	0.48	0.56	0.64	0.72	0.80	0.88	0.96	1.04	1.12	1.20	1.28	1.36	1.44	1.52	1.60	1.68	1.76	1.84	1.92	2.00	2.08	2.16	2.24	2.32	2.40
0.09	0.09	0.18	0.27	0.36	0.45	0.54	0.63	0.72	0.81	0.90	0.99	1.08	1.17	1.26	1.35	1.44	1.53	1.62	1.71	1.80	1.89	1.98	2.07	2.16	2.25	2.34	2.43	2.52	2.61	2.70
0.10	0.10	0.20	0.30	0.40	0.50	0.60	0.70	0.80	0.90	1.00	1.10	1.20	1.30	1.40	1.50	1.60	1.70	1.80	1.90	2.00	2.10	2.20	2.30	2.40	2.50	2.60	2.70	2.80	2.90	3.00
0.11	0.11	0.22	0.33	0.44	0.55	0.66	0.77	0.88	0.99	1.10	1.21	1.32	1.43	1.54	1.65	1.76	1.87	1.98	2.09	2.20	2.31	2.42	2.53	2.64	2.75	2.86	2.97	3.08	3.19	3.30
0.12	0.12	0.24	0.36	0.48	0.60	0.72	0.84	0.96	1.08	1.20	1.32	1.44	1.56	1.68	1.80	1.92	2.04	2.16	2.28	2.40	2.52	2.64	2.76	2.88	3.00	3.12	3.24	3.36	3.48	3.60
0.13	0.13	0.26	0.39	0.51	0.63	0.75	0.87	0.99	1.11	1.23	1.35	1.47	1.59	1.71	1.83	1.95	2.07	2.19	2.31	2.43	2.55	2.67	2.79	2.91	3.03	3.15	3.27	3.39	3.51	3.63
0.14	0.14	0.28	0.42	0.56	0.70	0.84	0.98	1.12	1.26	1.40	1.54	1.68	1.82	1.96	2.10	2.24	2.38	2.52	2.66	2.80	2.94	3.08	3.22	3.36	3.50	3.64	3.78	3.92	4.06	4.20
0.15	0.15	0.30	0.45	0.60	0.75	0.90	1.05	1.20	1.35	1.50	1.65	1.80	1.95	2.10	2.25	2.40	2.55	2.70	2.85	3.00	3.15	3.30	3.45	3.60	3.75	3.90	4.05	4.20	4.35	4.50
0.16	0.16	0.32	0.48	0.64	0.80	0.96	1.12	1.28	1.44	1.60	1.76	1.92	2.08	2.24	2.40	2.56	2.72	2.88	3.04	3.20	3.36	3.52	3.68	3.84	4.00	4.16	4.32	4.48	4.64	4.80
0.17	0.17	0.34	0.51	0.68	0.85	1.02	1.19	1.36	1.53	1.70	1.87	2.04	2.21	2.38	2.55	2.72	2.89	3.06	3.23	3.40	3.57	3.74	3.91	4.08	4.25	4.42	4.59	4.76	4.93	5.10
0.18	0.18	0.36	0.54	0.72	0.90	1.08	1.26	1.44	1.62	1.80	1.98	2.16	2.34	2.52	2.70	2.88	3.06	3.24	3.42	3.60	3.78	3.96	4.14	4.32	4.50	4.68	4.86	5.04	5.22	5.40
0.19	0.19	0.38	0.57	0.76	0.95	1.14	1.33	1.52	1.71	1.90	2.09	2.28	2.47	2.66	2.85	3.04	3.23	3.42	3.61	3.80	3.99	4.18	4.37	4.56	4.75	4.94	5.13	5.32	5.51	5.70
0.20	0.20	0.40	0.60	0.80	1.00	1.20	1.40	1.60	1.80	2.00	2.20	2.40	2.60	2.80	3.00	3.20	3.40	3.60	3.80	4.00	4.20	4.40	4.60	4.80	5.00	5.20	5.40	5.60	5.80	6.00
0.21	0.21	0.42	0.63	0.84	1.05	1.26	1.47	1.68	1.89	2.10	2.31	2.52	2.73	2.94	3.15	3.36	3.57	3.78	3.99	4.20	4.41	4.62	4.83	5.04	5.25	5.46	5.67	5.88	6.09	6.30
0.22	0.22	0.44	0.66	0.88	1.10	1.32	1.54	1.76	1.98	2.20	2.42	2.64	2.86	3.08	3.30	3.52	3.74	3.96	4.18	4.40	4.62	4.84	5.06	5.28	5.50	5.72	5.94	6.16	6.38	6.60
0.23	0.23	0.46	0.69	0.92	1.14	1.36	1.58	1.80	2.02	2.24	2.46	2.68	2.90	3.12	3.34	3.56	3.78	4.00	4.22	4.44	4.66	4.88	5.10	5.32	5.54	5.76	5.98	6.20	6.42	6.64
0.24	0.24	0.48	0.72	0.96	1.20	1.44	1.68	1.92	2.16	2.40	2.64	2.88	3.12	3.36	3.60	3.84	4.08	4.32	4.56	4.80	5.04	5.28	5.52	5.76	6.00	6.24	6.48	6.72	6.96	7.20
0.25	0.25	0.50	0.75	1.00	1.25	1.50	1.75	2.00	2.25	2.50	2.75	3.00	3.25	3.50	3.75	4.00	4.25	4.50	4.75	5.00	5.25	5.50	5.75	6.00	6.25	6.50	6.75	7.00	7.25	7.50
0.26	0.26	0.52	0.78	1.04	1.30	1.56	1.82	2.08	2.34	2.60	2.86	3.12	3.38	3.64	3.90	4.16	4.42	4.68	4.94	5.20	5.46	5.72	5.98	6.24	6.50	6.76	7.02	7.28	7.54	7.80
0.27	0.27	0.54	0.81	1.08	1.36	1.64	1.92	2.20	2.48	2.76	3.04	3.32	3.60	3.88	4.16	4.44	4.72	5.00	5.28	5.56	5.84	6.12	6.40	6.68	6.96	7.24	7.52	7.80	8.08	8.36
0.28	0.28	0.56	0.84	1.12	1.40	1.68	1.96	2.24	2.52	2.80	3.08	3.36	3.64	3.92	4.20	4.48	4.76	5.04	5.32	5.60	5.88	6.16	6.44	6.72	7.00	7.28	7.56	7.84	8.12	8.40
0.29	0.29	0.58	0.87	1.16	1.44	1.72	2.00	2.28	2.56	2.84	3.12	3.40	3.68	3.96	4.24	4.52	4.80	5.08	5.36	5.64	5.92	6.20	6.48	6.76	7.04	7.32	7.60	7.88	8.16	8.44
0.30	0.30	0.60	0.90	1.20	1.50	1.80	2.10	2.40	2.70	3.00	3.30	3.60	3.90	4.20	4.50	4.80	5.10	5.40	5.70	6.00	6.30	6.60	6.90	7.20	7.50	7.80	8.10	8.40	8.70	9.00

TABLE II (cont'd)

GENERATOR OUTPUT IN VOLTS AC FOR GAP LENGTH = 0.075 INCHES

REFC THICKNESS INCHES	AMPLITUDE IN INCHES	0.01	0.02	0.03	0.04	0.05	0.06	0.07	0.08	0.09	0.10	0.11	0.12	0.13	0.14	0.15	0.16	0.17	0.18	0.19	0.20	0.21	0.22	0.23	0.24	0.25	0.26	0.27	0.28	0.29
.013	3	7	10	13	17	20	24	27	31	35	39	43	47	51	54	58	62	66	70	74	77	81	85	89	93	97	101	105	109	113
.013	3	7	11	14	18	21	25	29	33	37	41	45	49	53	57	61	65	69	73	77	81	85	89	93	97	101	105	109	113	117
.014	4	7	11	15	19	23	27	31	35	39	43	47	51	55	59	63	67	71	75	79	83	87	91	95	99	103	107	111	115	119
.014	4	8	12	16	20	24	28	32	36	40	44	48	52	56	60	64	68	72	76	80	84	88	92	96	100	104	108	112	116	120
.014	4	8	13	17	21	25	29	33	37	41	45	49	53	57	61	65	69	73	77	81	85	89	93	97	101	105	109	113	117	121
.020	5	9	13	17	21	25	29	33	37	41	45	49	53	57	61	65	69	73	77	81	85	89	93	97	101	105	109	113	117	121
.022	5	9	14	18	22	26	30	34	38	42	46	50	54	58	62	66	70	74	78	82	86	90	94	98	102	106	110	114	118	122
.024	5	10	15	20	24	28	32	36	40	44	48	52	56	60	64	68	72	76	80	84	88	92	96	100	104	108	112	116	120	124
.024	5	11	16	22	28	34	40	46	52	58	64	70	76	82	88	94	100	106	112	118	124	130	136	142	148	154	160	166	172	178
.024	6	11	18	24	30	36	42	48	54	60	66	72	78	84	90	96	102	108	114	120	126	132	138	144	150	156	162	168	174	180
.030	6	13	19	26	32	38	44	50	56	62	68	74	80	86	92	98	104	110	116	122	128	134	140	146	152	158	164	170	176	182
.032	7	14	21	28	35	42	50	57	64	71	78	85	92	99	106	113	120	127	134	141	148	155	162	169	176	183	190	197	204	211
.034	7	15	22	30	38	46	54	62	70	78	86	94	102	110	118	126	134	142	150	158	166	174	182	190	198	206	214	222	230	238
.034	8	16	24	33	42	51	60	71	81	91	101	111	121	131	141	151	161	171	181	191	201	211	221	231	241	251	261	271	281	291
.034	9	17	26	36	46	56	66	76	86	96	106	116	126	136	146	156	166	176	186	196	206	216	226	236	246	256	266	276	286	296
.040	10	20	30	40	51	62	74	85	96	107	118	129	140	151	162	173	184	195	206	217	228	239	250	261	272	283	294	305	316	327
.042	11	22	33	45	57	70	83	95	107	119	131	143	155	167	179	191	203	215	227	239	251	263	275	287	299	311	323	335	347	359
.042	12	24	37	50	64	78	92	106	120	134	148	162	176	190	204	218	232	246	260	274	288	302	316	330	344	358	372	386	400	414
.044	14	27	42	57	72	87	102	117	132	147	162	177	192	207	222	237	252	267	282	297	312	327	342	357	372	387	402	417	432	447
.044	15	31	48	65	83	101	119	137	155	173	191	209	227	245	263	281	299	317	335	353	371	389	407	425	443	461	479	497	515	533
.050	16	34	55	75	97	121	145	169	193	217	241	265	289	313	337	361	385	409	433	457	481	505	529	553	577	601	625	649	673	697
.052	21	42	64	88	114	140	166	192	218	244	270	296	322	348	374	400	426	452	478	504	530	556	582	608	634	660	686	712	738	764
.054	24	49	74	103	133	163	193	223	253	283	313	343	373	403	433	463	493	523	553	583	613	643	673	703	733	763	793	823	853	883
.056	29	60	92	124	157	190	223	256	289	322	355	388	421	454	487	520	553	586	619	652	685	718	751	784	817	850	883	916	949	982
.058	34	73	114	161	210	259	308	357	406	455	504	553	602	651	700	749	798	847	896	945	994	1043	1092	1141	1190	1239	1288	1337	1386	1435
.060	45	93	146	210	290	384	478	572	666	760	854	948	1042	1136	1230	1324	1418	1512	1606	1700	1794	1888	1982	2076	2170	2264	2358	2452	2546	2640
.062	49	124	186	268	362	456	550	644	738	832	926	1020	1114	1208	1302	1396	1490	1584	1678	1772	1866	1960	2054	2148	2242	2336	2430	2524	2618	2712
.064	51	170	240	331	435	539	643	747	851	955	1059	1163	1267	1371	1475	1579	1683	1787	1891	1995	2099	2203	2307	2411	2515	2619	2723	2827	2931	3035
.066	119	256	400	566	742	928	1114	1300	1486	1672	1858	2044	2230	2416	2602	2788	2974	3160	3346	3532	3718	3904	4090	4276	4462	4648	4834	5020	5206	5392
.068	103	435	680	935	1190	1445	1700	1955	2210	2465	2720	2975	3230	3485	3740	3995	4250	4505	4760	5015	5270	5525	5780	6035	6290	6545	6800	7055	7310	7565
.070	375	760	1145	1530	1915	2300	2685	3070	3455	3840	4225	4610	4995	5380	5765	6150	6535	6920	7305	7690	8075	8460	8845	9230	9615	10000	10385	10770	11155	11540

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